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Brush tyre models for large camber angles and steering speeds

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ABSTRACT

In this paper, we discuss three improved brush models. The first one deals with the coupling between the slip and spin parameters and is valid for relatively high steering speed and small camber angles; the second one is more complex and considers the presence of a two-dimensional velocity field inside the contact patch due to large camber angles; the third one is more general and combines both the previous formulations. For the last two models, the investigation is conducted with respect to a rectangular contact patch, for which we show that three different regions can be identified, each of them corresponding to a different steady-state solution for the deflection of the bristle. Furthermore, from the transient analysis it emerges that each region can be in turn separated into an area in which steady-state conditions reign and another one in which the transient solution takes place. An asymptotic analysis is carried out for the three models and it is shown that the solutions are equivalent to the ones predicted by the standard brush theory for small values of the spin ratio and camber angle. Finally, a comparison is performed amongst the models to highlight the differences in the predicted tyre characteristics.

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Tyre modelling; brush models; nonlinear theory; two-dimensional theory; transient dynamics; transport equation

Acronyms and abbreviations

In the present paper, we use some abbreviations. We summarise here the main acronyms used throughout the paper:

Acronym	Explanation
1DM	One-dimensional brush model
1DCM	One-dimensional brush model for coupled slips and spin
2DM	Two-dimensional brush model
2DCM	Two-dimensional brush model for coupled slips and spin
BC	Boundary condition
IC	Initial condition
ID	Initial data
IFT	Implicit Function Theorem

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ODE Ordinary differential equation
PDE Partial differential equation

1. Introduction

Tyres constitute the primary interface between ground vehicles and the environment. Their peculiar behavioural properties allow them to undergo deformations which cause longitudinal and lateral forces to arise inside the contact patch. Thus, a reliable prediction of tyre characteristics is crucial when it comes to real-time optimisation of the vehicle overall performance. Another critical aspect concerns tyre characterisation. Indeed, tyre data can be often collected only under highly controlled conditions, which might be difficult to replicate with precision due to the long and expensive procedures. Also, these tests are mainly carried out on very smooth surfaces which do not reflect the real working conditions.

Accurate tyre modelling has gained a vital importance during the last five decades, and multiple approaches have been proposed depending on the specific application. In particular, the most famous model is the so-called Pacejka's Magic Formula (MF): a wide set of empirical equations which are able to fit tyre characteristics with astonishing precision for different working ranges [1,2]. The MF has been perfected over the years to cover several cases which were not contemplated by the earlier versions, being today the ultimate tyre model. For example, a great deal of research was devoted by both Besselink and Pacejka [3,4] to account for the effect of the inflation pressure and for transients. More recently, Farroni [5–8] integrated the MF description with a full-physical thermal model in order to include the contribution of the tyre temperature. As a result, the MF coefficients have been redefined as temperature-dependent functions. However, the main drawback of MF formulation is that it relies on a large number of fitting parameters, whose physical meaning is not always easy to interpret.

An antipodal description, full-physical tyre models – like FTire[®] [9–11] and CDTire [12] – are instead capable of an accurate description of the local phenomena occurring inside the contact patch. In particular, the bordering on perfection FTire[®] combines multibody dynamics features with nonlinear curved beam equations and appropriate circumferential extensibility of the tyre, and represents today a standard in the context of advanced driving simulations. Indeed, in spite of their complexities, these models are able to run in real-time. However, simple analytical investigations are merely impossible to conduct because of the extremely detailed modelling.

Other physical theories are based on an exhaustive description of the frictional phenomena which take place in the contact patch. Some of them have been derived from more generalising models employed in the context of tribological studies. Perhaps, the most famous friction-based tyre model is the LuGre [13], initially developed during a collaboration between the Lund and Grenoble universities and then extended by several authors towards disparate applications, including vehicle dynamics and other automotive-related fields [14–20]. In particular, in the LuGre formulation, the planar contact stresses exerted at the tyre-road interface and the friction-induced hysteresis are described in terms of an inertial friction state which can be identified with the deformation of a material point of the tyre tread contacting the ground. The governing equation of the model is a one-dimensional inhomogeneous transport equation, which is usually discretised in space and then solved in time.

On the other hand, brush models [1,2,21–24] represent a more naive but still insightful approach when it comes to a pure physical description of the tyre. They are derived from a simplified version of the classic theory of contact mechanics [25–31], are grounded on few assumptions, and are still accurate enough to reflect the main variations in tyre characteristics due to different operative conditions. Furthermore, the reduced order in the number of required parameters allows for essential speculations about the governing physics of the tyre-road kinematics, providing a deep understanding of the basic phenomena occurring in the contact patch. The underlying idea is that the tyre-wheel system can be assumed to be a rigid body provided with bristles which only deform in longitudinal and lateral direction inside the contact patch. The governing equations are hence expressed in Eulerian form.

In particular, the brush theory was established as main mathematical foundation for tyre modelling since the earliest studies conducted by Pacejka on shimmy phenomena. Indeed, he used the brush theory to investigate the local deformation of string-like models during lateral manoeuvres [32]. Pacejka found that the displacement field of the tyre carcass can be deduced starting from the tyre-road kinematic equations and managed to quantify the generalised forces acting on the wheel hub in the frequency domain. The importance of the brush models has been then consolidated and sometimes revamped over the years thanks to several famous works relating to different fields. For example, Higuchi [33,34] modelled the effect of large camber angles and derived an alternative version of the tyre-road kinematic equations. Takacs [35,36] continued with very remarkable and insightful investigations about shimmy and micro-shimmy related phenomena [37–41], combining brush models with more advanced approaches. Svendenius [42–44] investigated the effect of more realistic shapes of the vertical pressure inside the contact patch and developed a semi-empirical formulation to easily account for the contribution of non-negligible camber angles. Together with other authors [45–47], he used the brush models for the purpose of real-time friction estimation. In his book, Guiggiani [21] redefined the classic science of vehicle dynamics providing a more mathematically rigorous and systematic approach. From a tyre modelling perspective, he emphasised the crucial role played by the brush models because of their physical nature, and analysed the connected transient problem in a very general case. Other authors have studied unsteady-state conditions by using numerical methods [48,49] or proposed enhanced formulations to extend the classic theory towards complex geometries [50,51].

Recently, the authors have developed an exhaustive theory which captures the dynamics of the bristles in the contact patch [52]. A class of simplified formulae has also been introduced to deal with more complex scenarios and take into account the major effects connected with the tyre carcass dynamics.

In the present paper, we extend our previous investigation to analyse a wide class of theoretically unexplored problems. More specifically, we first generalise the classic boundary conditions prescribed by the standard brush theory by expressing them in terms of the inflowing flux. Then, we derive three different models in order of complexity. The first one deals with sufficiently high steering speeds to excite the coupling between the longitudinal and lateral deflections of the bristle and the slip and spin parameters, respectively. This variant of the classic brush model predicts a correlation between the longitudinal deflection and the lateral slip and *vice versa*, and can be used to reflect working conditions for which the cornering plays a crucial role but the camber angle is instead

limited. We refer to this theory as the *one-dimensional model for coupled slips and spin* (1DCM).

The second model – renamed *two-dimensional model* (2DM) – is valid for arbitrary large values of the camber angle and deals with the presence of a two-dimensional velocity field inside the contact patch. In this case, the governing equations result in a more complicated formulation. For a rectangular contact patch, we prove the existence of three different regions in which the steady-state solution is given by a different analytical expression. For both the above mentioned models we perform an asymptotic analysis to ascertain that they are equivalent to the classic one for small values of the spin parameter and camber angle, respectively. The last model is the complete one, referred to as *two-dimensional model for coupled slips and spin* (2DCM), and accounts for both the presence of the coupling between the slips and spin parameters and for a two-dimensional velocity field. Also in this case, we conduct a complete analysis in the case of a rectangular contact shape. The resulting expressions for the deflections of the bristles inside the contact patch show similar features with the ones predicted by both the two previous theories. Since the characteristics equations are the same as for the 2DM, we can also identify the same regions in the contact patch in which the steady-state and transient solution for the bristle displacement assumes a different analytical expression. Finally, an asymptotic analysis is carried out to verify that the 2DCM is consistent with the 2DM for sufficiently small values of the steering speed.

We stress that the models presented in work paper are not intended as an alternative to the classic brush theory, which is sufficient in most of the cases and much easier to implement in simulations. Instead, they can be thought as an effective theoretical foundation to investigate analytically some phenomena which are often studied by means of advanced FEM or multibody models.

This paper is organised as follows: in Section 1, the tyre-road contact equations are presented in the most general form. The definition of the generalised slip quantities is introduced and the main variables and assumptions are explained. In Section 2, the 1DCM is presented. The governing equations are solved for both the steady-state and the transient cases and some comparisons are performed with the results predicted by the classic brush theory. In Section 3, the 2DM is derived and we show some results concerning the shape of the solution for the steady-state problem. The transient solution is also derived and the asymptotic analysis is carried out. In Section 4, the 2DCM is discussed and an analytical solution for the deflection of the bristles in each subdomain of the contact patch is given. Consistency with the 2DM is verified for the longitudinal deflection in the main area of the contact patch. Finally, the main conclusions are drawn in Section 6 and further steps in research are outlined.

2. Tyre-road contact mechanics equations

Let us consider a reference frame $Oxyz$ with unit vectors $(\hat{e}_x, \hat{e}_y, \hat{e}_z)$, whose origin O is located in the centroid of the *contact patch*; the axes are oriented according to the SAE system: the x axis is directed towards the longitudinal direction of motion, the z axis points downward and the y axis lies in on the road surface and is oriented so that the coordinate system is right-handed. The contact patch is defined mathematically as a closed set \mathcal{P} , whose interior and boundary are denoted with $\overset{\circ}{\mathcal{P}}$ and $\partial \mathcal{P}$, respectively. The contact patch

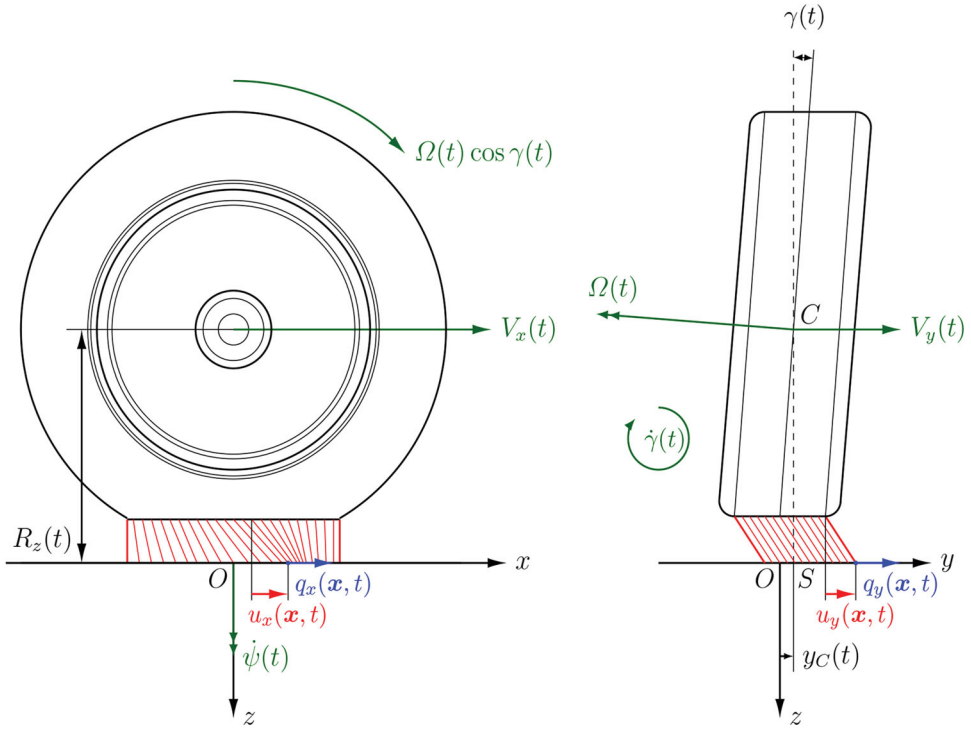


Figure 1. Tyre-road schematics in the Oxz and Oyz planes (left and right-hand side figures, respectively). The geometric parameters are drawn in black, the kinematic quantities in green, the generalised forces in blue and the elastic displacements in red. The dimension and the deformation of the bristles have been exaggerated to facilitate the visualisation.

collects all the points $\mathbf{x} \in \mathbb{R}^3$ of the tyre which make contact with the road (represented by the Oxy plane). The road is modelled as a perfect homogeneous, isotropic flat surface, without any irregularity.

During the rolling of the tyre, a quantity can evolve over the time $t \in \mathbb{R}_{\geq 0}$ or, equivalently, over a peculiar distance $s \in \mathbb{R}_{\geq 0}$ which we call *travelled* or *rolling distance*; its meaning will be clarified in the following. In particular, at each point $\mathbf{x} \in \mathcal{P}$ we associate a three-dimensional speed field $\dot{\mathbf{x}} = \mathbf{v}(\mathbf{x}, t) = v_x(\mathbf{x}, t)\hat{\mathbf{e}}_x + v_y(\mathbf{x}, t)\hat{\mathbf{e}}_y + v_z(\mathbf{x}, t)\hat{\mathbf{e}}_z$ and a finite vector displacement $\mathbf{u}(\mathbf{x}, t) = u_x(\mathbf{x}, t)\hat{\mathbf{e}}_x + u_y(\mathbf{x}, t)\hat{\mathbf{e}}_y + u_z(\mathbf{x}, t)\hat{\mathbf{e}}_z$, which represents the relative deformation of the material point located at the coordinate \mathbf{x} with respect to its initial configuration. In the brush model, the deformation of a material point is also interpreted as the deformation of a *bristle* attached to the tyre; hence, we refer sometimes to the deformation, deflection or displacement of a bristle. Each material point may be also subjected to a force per unit of area, which may vary both in space and time, $\mathbf{q}(\mathbf{x}, t) = q_x(\mathbf{x}, t)\hat{\mathbf{e}}_x + q_y(\mathbf{x}, t)\hat{\mathbf{e}}_y + q_z(\mathbf{x}, t)\hat{\mathbf{e}}_z$. The whole situation is illustrated in Figure 1.

Since we only deal with the planar problem, we consider the tangential (or planar¹) and normal vectors to the road plane, denoted with $\mathbf{t} = [x \ y \ 0]^T$ and $\mathbf{n} = [0 \ 0 \ z]^T$. Accordingly, we also define the tangential velocity field, the tangential displacement vectors and the tangential stress vector as $\mathbf{v}_t = v_x(\mathbf{x}, t)\hat{\mathbf{e}}_x + v_y(\mathbf{x}, t)\hat{\mathbf{e}}_y$, $\mathbf{u}_t = u_x(\mathbf{x}, t)\hat{\mathbf{e}}_x +$

$u_y(\mathbf{x}, t)\hat{\mathbf{e}}_y$ and $\mathbf{q}_t(\mathbf{x}, t) = q_x(\mathbf{x}, t)\hat{\mathbf{e}}_x + q_y(\mathbf{x}, t)\hat{\mathbf{e}}_y$, so that $\mathbf{v}(\mathbf{x}, t) = [\mathbf{v}_t(\mathbf{x}, t) \ v_z(\mathbf{x}, t)]^T$, $\mathbf{u}(\mathbf{x}, t) = [\mathbf{u}_t(\mathbf{x}, t) \ u_z(\mathbf{x}, t)]^T$ and $\mathbf{q}(\mathbf{x}, t) = [\mathbf{q}_t(\mathbf{x}, t) \ q_z(\mathbf{x}, t)]^T$. Sometimes, in this paper, we also call $\mathbf{q}_t(\mathbf{x}, t)$ *shear stress vector*.

Finally, we assume that the planar problem can be decoupled from the vertical one; this implies that the vertical component of each quantity defined so far is not influenced by the tangential ones. In particular, we assume that no modification occurs in the original vertical pressure distribution $q_z(\mathbf{x}, t)$ due to friction-related phenomena.

The relative speed between a material point inside the contact patch and the road is called *micro-sliding speed* indicated with $\mathbf{v}_s(\mathbf{x}, t)$ ³ and the tangential micro-sliding speed is defined pertinently as $\mathbf{v}_{st}(\mathbf{x}, t)$. Adopting a simple Coulomb friction model, the fundamental equations governing the tyre-road contact mechanics can be now formulated as follows [21]:

$$\mathbf{v}_{st}(\mathbf{x}, t) = \mathbf{0} \iff q_t(\mathbf{x}, t) \leq \mu_s q_z(\mathbf{x}, t), \quad (1a)$$

$$\mathbf{q}_t(\mathbf{x}, t) = -\mu_d q_z(\mathbf{x}, t) \frac{\mathbf{v}_{st}(\mathbf{x}, t)}{\|\mathbf{v}_{st}(\mathbf{x}, t)\|} \iff \mathbf{v}_{st}(\mathbf{x}, t) \neq \mathbf{0}, \quad (1b)$$

where $q_t(\mathbf{x}, t) = \|\mathbf{q}_t(\mathbf{x}, t)\|$ and μ_s and μ_d are the static and dynamic friction coefficient, respectively. In general, both can be made dependent explicitly on the vector position \mathbf{x} or on the sliding speed $\mathbf{v}_t(\mathbf{x}, t)$; however, to keep the complexity of the analysis within acceptable levels, here we only use constant values for μ_s and μ_d . For more sophisticated formulations and extensive discussion, the reader may instead refer to [53–55]. There are, of course some limitations connected with the latter simplification. More specifically, Equation (1) are only valid under the assumption of memoryless friction, and hence dynamic behaviours due to complex phenomena, e.g. viscoelasticity, are automatically neglected.

To solve the above Equation (1), two other sets of relations are needed: the *tyre-road kinematic equations* and the *constitutive relations*. The first set prescribes a relation between the sliding speed and the deformation of the tyre inside the contact patch; the latter the relation between the aforementioned deformation and the tangential stress acting on each material point.

2.1. Tyre-road kinematic equations

In the present paper, we assume that – from a frictional perspective – there is only one sticking zone (in which the bristles are in adhesion condition) and one sliding zone (in which finite sliding occurs between the bristles and the road, i.e. $\mathbf{v}_{st}(\mathbf{x}, t) \neq \mathbf{0}$) inside the open set $\mathring{\mathcal{P}}$. This is a strong assumption, but leads to an enormous simplification of the problem. Indeed, in this way, the domain in which the partial differential equations (PDEs) ruling the tyre-road kinematics are defined coincides with the interior of the contact patch $\mathring{\mathcal{P}}$, and the corresponding BCs can be formulated uniquely on $\partial \mathring{\mathcal{P}}$. More exhaustive models of friction, which do not rely on the latter simplifying hypothesis, can be found in [16–20, 56, 57]. It is also worth pointing out that some efforts have been devoted to remove the assumption of a single sticking zone also in the context of the brush theory, for example recently by Guiggiani [21], but the investigation becomes feasible only numerically by means of

iterative algorithms. Here we want to conduct an investigation to derive, where possible, a closed-form expression for the tangential deformation $\mathbf{u}_t(\mathbf{x}, t)$ of the bristle, and hence we try to simplify the analysis as much as possible.

The micro-sliding speed $\mathbf{v}_{st}(\mathbf{x}, t)$ reads, in vector form, Pacejka [2], Guiggiani [21], Limebeer and Massaro [23], Romano et al. [24], Romano et al. [52]

$$\begin{aligned} \mathbf{v}_{st}(\mathbf{x}, t) = & \mathbf{V}_{st}(t) + \dot{\mathbf{x}}_C(t) + \mathbf{W}_z(t)(\mathbf{x} - \mathbf{x}_C(t) + \varepsilon_\psi(t)\mathbf{u}_t(\mathbf{x}, t)) \\ & + \frac{d\mathbf{u}_t(\mathbf{x}, t)}{dt}, \quad \mathbf{x} \in \mathcal{P}, \quad t \in \mathbb{R}_{>0}, \end{aligned} \quad (2)$$

with

$$\mathbf{V}_{st}(t) \triangleq \begin{bmatrix} V_{sx}(t) & V_{sy}(t) \end{bmatrix}^T = \begin{bmatrix} V_x(t) - \Omega(t)R_r(t) & V_y(t) - \dot{\gamma}(t)R_z(t) \end{bmatrix}^T, \quad (3a)$$

$$\mathbf{W}_z(t) \triangleq \begin{bmatrix} 0 & -\omega_z(t) \\ \omega_z(t) & 0 \end{bmatrix}, \quad (3b)$$

$$\omega_z(t) \triangleq \dot{\psi}(t) - \Omega(t) \sin \gamma(t), \quad (3c)$$

$$\mathbf{u}_t(\mathbf{x}, t) \triangleq \begin{bmatrix} u_x(\mathbf{x}, t) & u_y(\mathbf{x}, t) \end{bmatrix}^T, \quad (3d)$$

$$\varepsilon_\gamma(t) \triangleq -\frac{\Omega(t) \sin \gamma(t)}{\omega_z(t)}, \quad (3e)$$

$$\varepsilon_\psi(t) \triangleq 1 - \varepsilon_\gamma(t) = \frac{\dot{\psi}(t)}{\omega_z(t)}, \quad (3f)$$

where $\mathbf{V}_{st}(t) \triangleq \mathbf{V}_{St}(t) = [V_{Sx}(t) \ V_{Sy}(t)]^T$ is the *macro-sliding velocity*⁴ and coincides with the velocity of a hypothetical point S representing the projection of the centre C of the wheel hub onto the road plane, $V_x(t) \triangleq V_{Cx}(t)$ and $V_y(t) \triangleq V_{Cy}(t)$ are the longitudinal and lateral speed of the wheel hub, whose planar coordinates are given by the vector $\mathbf{x}_C = x_C \hat{\mathbf{e}}_x + y_C \hat{\mathbf{e}}_y$, $\Omega(t)$ is the rolling speed of the tyre around its axis, $R_r(t) = R_z(t) \cos \gamma(t)$ is the so-called *rolling radius*, $\gamma(t)$ is the camber angle and $\dot{\psi}(t)$ is the steering angular speed. The geometrical radius $-R_z(t)\hat{\mathbf{e}}_x$ denotes the vertical coordinate of the centre of the wheel hub with respect to the reference frame $Oxyz$. Finally, in the authors' knowledge, the quantities $\varepsilon_\gamma(t)$ and $\varepsilon_\psi(t)$ have never been defined in literature; we call them *camber* and *steering ratio*, respectively. They represent the contribution of the two components of the angular speed in z direction (the first one due to the camber angle, the second to the steering speed) to the total value of $\omega_z(t)$.

It is crucial to understand that, since we are analysing the tyre-road interaction by means of the Eulerian approach, the total derivative of a quantity is given by $d/dt = \partial/\partial t + \mathbf{v}(\mathbf{x}, t) \cdot \nabla$. However, since we are only focussing on the tangential problem, $\partial/\partial z = 0$ and hence we may instead use $d/dt = \partial/\partial t + \mathbf{v}_t(\mathbf{x}, t) \cdot \nabla_t$, where ∇_t collects the tangential components of the gradient.

With the premises above, Equation (2) can be recast in scalar form as

$$v_{sx}(\mathbf{x}, t) = V_{sx}(t) + \dot{x}_C(t) - \omega_z(t)(y - y_C(t) + \varepsilon_\psi(t)u_y(\mathbf{x}, t)) + \frac{du_x(\mathbf{x}, t)}{dt}, \quad \mathbf{x} \in \overset{\circ}{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (4a)$$

$$v_{sy}(\mathbf{x}, t) = V_{sy}(t) + \dot{y}_C(t) + \omega_z(t)(x - x_C(t) + \varepsilon_\psi(t)u_x(\mathbf{x}, t)) + \frac{du_y(\mathbf{x}, t)}{dt}, \quad \mathbf{x} \in \overset{\circ}{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (4b)$$

with

$$\frac{du_x(\mathbf{x}, t)}{dt} = \frac{\partial u_x(\mathbf{x}, t)}{\partial t} + v_x(\mathbf{x}, t) \frac{\partial u_x(\mathbf{x}, t)}{\partial x} + v_y(\mathbf{x}, t) \frac{\partial u_x(\mathbf{x}, t)}{\partial y}, \quad (5a)$$

$$\frac{du_y(\mathbf{x}, t)}{dt} = \frac{\partial u_y(\mathbf{x}, t)}{\partial t} + v_x(\mathbf{x}, t) \frac{\partial u_y(\mathbf{x}, t)}{\partial x} + v_y(\mathbf{x}, t) \frac{\partial u_y(\mathbf{x}, t)}{\partial y}. \quad (5b)$$

In the classic brush theory, which deals with a one-dimensional velocity field $\mathbf{v}_t(\mathbf{x}, t) \approx -V_r(t)\hat{\mathbf{e}}_x$, the coupling between the longitudinal and tangential displacement of the bristle is also neglected, i.e. $\varepsilon_\psi = 0$. We refer to the classic theory simply as *one-dimensional theory* (IDM). Instead, in the following analysis, we will make different assumptions on the components $v_x(\mathbf{x}, t)$ and $v_y(\mathbf{x}, t)$ of the velocity field, as well as on the parameter $\varepsilon_\psi(t)$, to properly account for different operative conditions of the tyre. The boundary conditions (BCs) for the problem under consideration can be stated inherently depending on the specific structure of the velocity field.

2.1.1. Slip parameters

The solution of the tyre-road kinematic equations simplify if the analysis is conducted with respect to nondimensional quantities, called *slip variables* or *parameters*.⁵ In this way, the problem can be made independent of the rolling speed of the tyre. Thus, we introduce the *theoretical slip parameters* reading

$$\boldsymbol{\sigma}(t) = \begin{bmatrix} \sigma_x(t) \\ \sigma_y(t) \end{bmatrix} \triangleq -\frac{\mathbf{V}_{st}(t)}{V_r(t)}, \quad (6a)$$

$$\varphi(t) \triangleq -\frac{\omega_z(t)}{\Omega(t)}, \quad (6b)$$

where the quantity $V_r(t) \triangleq \Omega(t)R_r(t)$ is referred to as *rolling speed* and the subscript *r* stands for *rolling*. The quantities $\sigma_x(t)$ and $\sigma_y(t)$ are called *longitudinal* and *lateral slip*, respectively, whilst φ is referred to both as *rotational slip* or *spin*. We point out that, unlikely many authors which define the spin as $\varphi \triangleq -\omega_z(t)/V_r(t)$, we have chosen to divide simply by $\Omega(t)$; in this way, $\boldsymbol{\sigma}(t)$ and $\varphi(t)$ are both nondimensional quantities. For some applications, it may also be useful to define the *transient slip* as in [21], which incorporates in its definition the transient deflection of the tyre carcass. Generally speaking, in the following investigations we will neglect the carcass dynamics, even though it heavily affects the transient trend of the forces exerted at the tyre-road interface; this kind of study, however, is out of the scope of the present paper.

Finally, we introduce mathematically the notion of travelled distance as follows:

$$s \triangleq \int_0^t V_r(t') dt'. \quad (7)$$

In many cases, we will use the travelled distance as independent coordinate in place of the time variable.

2.2. Constitutive relation

In spite of the viscoelastic nature of the tyre, for sake of simplicity it has been commonly established in literature to assume linear elasticity [21], i.e. a constitutive relation of the type

$$\mathbf{q}_t(\mathbf{x}, t) = \mathbf{K}_t \mathbf{u}_t(\mathbf{x}, t), \quad (8)$$

where the tangential stiffness matrix \mathbf{K}_t is defined as [21]

$$\mathbf{K}_t = \begin{bmatrix} k_x & 0 \\ 0 & k_y \end{bmatrix}. \quad (9)$$

2.3. Boundary and initial conditions

Equation (4) are two coupled PDEs – more specifically, *linear transport equations* – defined on a finite open domain $\mathring{\mathcal{P}}$. Thus, to guarantee the uniqueness of the solution, we need to prescribe a BC and an initial condition (IC). To formalise the BCs correctly, it is firstly necessary to define the *leading edge* \mathcal{L} , the *neutral edge* \mathcal{N} and the *trailing edge* \mathcal{T} [58] as

$$\mathcal{L} \triangleq \{ \mathbf{x} \in \partial \mathcal{P} \mid \mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x}) < 0 \}, \quad (10a)$$

$$\mathcal{N} \triangleq \{ \mathbf{x} \in \partial \mathcal{P} \mid \mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x}) = 0 \}, \quad (10b)$$

$$\mathcal{T} \triangleq \{ \mathbf{x} \in \partial \mathcal{P} \mid \mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x}) > 0 \}, \quad (10c)$$

where the unit vector $\hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x})$ represents the outer-pointing unit normal which lies in the $= Oxy$ plane. Note that the scalar product $\mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x})$ represents the elementary flux $d\Phi_{\partial \mathcal{P}}(\mathbf{v}_t(\mathbf{x})) / d\mathbf{x}$ through the boundary $\partial \mathcal{P}$ of the contact patch.

Indeed, the classic brush theory prescribes the continuity of the shear stress at the interface between the free portion of the half space and the interior of contact area. If a pure elastic constitutive relation is assumed, the direct consequence of is that the bristles inflowing into the contact patch must enter undeformed.

The previous boundary can be stated in mathematical terms as

$$\text{BC: } \mathbf{q}_t(\mathbf{x}, t) = \mathbf{K}_t \mathbf{u}_t(\mathbf{x}, t) = \mathbf{0} \iff \mathbf{u}_t(\mathbf{x}, t) = \mathbf{0}, \quad \mathbf{x} \in \mathcal{L}, \quad t \in \mathbb{R}_{>0}. \quad (11)$$

Basically, the previous relation imposes that the bristles must enter the contact patch undeformed, since the points $\mathbf{x} \in \mathcal{L}$ are the points inflowing into the contact patch \mathcal{P} .

Furthermore, we can partition the leading edge as follows

$$\mathcal{L} = \bigcup_{i \in \mathcal{I}} \mathcal{L}_i, \quad (12)$$

so that, with the necessary precautions needed to apply the *Implicit Function Theorem* (IFT), we can find two parametrisations for each subset \mathcal{L}_i , namely $x_{\mathcal{L}_i}(y)$ and $y_{\mathcal{L}_i}(x)$, such that we can write

$$\mathcal{L}_i = x_{\mathcal{L}_i}(y), \quad y_{i1} \leq y < y_{i2}, \quad z = 0, \quad (13a)$$

$$\mathcal{L}_i = y_{\mathcal{L}_i}(x), \quad x_{i1} \leq x < x_{i2}, \quad z = 0, \quad (13b)$$

the BC (11) can be recast as

$$\text{BC} : \mathbf{u}_t(x_{\mathcal{L}_i}(y), y, t) = \mathbf{0}, \quad \mathbf{x} \in \mathcal{L}_i, \quad t \in \mathbb{R}_{>0}, \quad (14a)$$

$$\text{BC} : \mathbf{u}_t(x, y_{\mathcal{L}_i}(x), t) = \mathbf{0}, \quad \mathbf{x} \in \mathcal{L}_i, \quad t \in \mathbb{R}_{>0}. \quad (14b)$$

Note that the above boundaries with the condition (10a) ensure that the boundary data (BDs) are never characteristic (the reader may refer, for example, to [59,60]).

As far as the IC, is concerned, we introduce the additional assumption of *vanishing sliding* [2], which corresponds to have infinite friction available inside the contact patch, i.e. $\mu_s \rightarrow \infty$. By doing this, we can formulate the IC on the whole interior $\overset{\circ}{\mathcal{P}}$ of the domain \mathcal{P} ⁶:

$$\text{IC} : \mathbf{u}_t(\mathbf{x}, 0) = \mathbf{u}_{t0}(\mathbf{x}), \quad \mathbf{x} \in \overset{\circ}{\mathcal{P}}, \quad (15)$$

for any $\mathbf{u}_{t0}(\mathbf{x}) \in C^1(\overset{\circ}{\mathcal{P}})$. The assumption of vanishing sliding will be eventually removed in Section 5. Another constraint which we introduce is that $\mathbf{u}_{t0}(\mathbf{x})|_{\mathcal{L}} = \mathbf{0}$. This can be justified by frictional considerations.

3. One-dimensional model for large steering speeds

In this section, we introduce a model which can be applied when the spin parameter is large. This case is relevant when the contribution of the steering speed is non-negligible, but the camber angle is sufficiently small to be disregarded ($\varepsilon_\gamma \approx 0$). We refer to this model as the *one-dimensional model for coupled slip and spin* (1DCM). The reason for such a definition is that, as it will be highlighted in the following, large spin values excite a coupling between the slip and spin parameters. Another assumption which we introduce is that all the geometric quantities, the speeds and the slip parameters in Equations (3) and (6) are constant over the time. Owing the premises above, the tangential velocity field can be approximated by the rolling speed as $dx/dt \approx -V_r$, so that

$$\begin{aligned} v_{sx}(\mathbf{x}, t) &= V_{sx} - \omega_z(y - y_C + \varepsilon_\psi u_y(\mathbf{x}, t)) + \frac{\partial u_x(\mathbf{x}, t)}{\partial t} \\ &\quad - V_r \frac{\partial u_x(\mathbf{x}, t)}{\partial x}, \quad \mathbf{x} \in \overset{\circ}{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \end{aligned} \quad (16a)$$

$$\begin{aligned} v_{sy}(\mathbf{x}, t) &= V_{sy} + \omega_z(x - x_C + \varepsilon_\psi u_x(\mathbf{x}, t)) + \frac{\partial u_y(\mathbf{x}, t)}{\partial t} \\ &\quad - V_r \frac{\partial u_y(\mathbf{x}, t)}{\partial x}, \quad \mathbf{x} \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}. \end{aligned} \quad (16b)$$

In the regions of the contact patch where no sliding occurs (i.e. $v_{st}(\mathbf{x}, t) = 0$), we can recast Equation (16) as

$$\frac{\partial u_x(\xi, s)}{\partial s} + \frac{\partial u_x(\xi, s)}{\partial \xi} = \sigma_x - \frac{\varphi}{R_r}(\eta - y_C + \varepsilon_\psi u_y(\xi, s)), \quad \xi \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (17a)$$

$$\frac{\partial u_y(\xi, s)}{\partial s} + \frac{\partial u_y(\xi, s)}{\partial \xi} = \sigma_y + \frac{\varphi}{R_r}(x_{\mathcal{L}}(\eta) - x_C - \xi + \varepsilon_\psi u_x(\xi, s)), \quad \xi \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (17b)$$

where we have introduced the new coordinate system

$$\xi = \begin{bmatrix} \xi \\ \eta \\ \zeta \end{bmatrix} \triangleq \begin{bmatrix} x_{\mathcal{L}}(y) - x \\ y \\ z \end{bmatrix}, \quad (18)$$

and replaced the time variable with the travelled distance $s = V_r t$. The variable ξ is often referred to as *distance from the entrance* and represents to the longitudinal coordinate measured from the leading edge. Indeed, the quantity $x_{\mathcal{L}}(y) = x_{\mathcal{L}}(\eta)$ represents the position of the leading coordinate at the lateral position $y(\eta)$.

Since, for the case under consideration, the velocity field is one-dimensional, the BC (14) and the IC (15) can be reformulated respectively as follows: black

$$\text{BC: } \mathbf{u}_t(0, \eta, s) = \mathbf{0}, \quad s \in \mathbb{R}_{>0}, \quad (19)$$

$$\text{IC: } \mathbf{u}_t(\xi, 0) = \mathbf{u}_{t0}(\xi), \quad \xi \in \mathring{\mathcal{P}}. \quad (20)$$

Solving (17a) for $u_y(\mathbf{x}, s)$ and substituting into (17b) yields

$$u_y(\xi, s) = \frac{1}{\varepsilon_\psi} \left[\frac{R_r}{\varphi} \left(\sigma_x - \frac{\partial u_x(\xi, s)}{\partial s} - \frac{\partial u_x(\xi, s)}{\partial \xi} \right) - \eta + y_C \right], \quad (21a)$$

$$\begin{aligned} &\frac{\partial^2 u_x(\xi, s)}{\partial s^2} + 2 \frac{\partial^2 u_x(\xi, s)}{\partial \xi \partial s} + \frac{\partial^2 u_x(\xi, s)}{\partial \xi^2} + \varepsilon_\psi^2 \frac{\varphi^2}{R_r^2} u_x(\xi, s) \\ &= -\varepsilon_\psi \frac{\varphi}{R_r} \left[\sigma_y + \frac{\varphi}{R_r} (x_{\mathcal{L}}(\eta) - x_C - \xi) \right]. \end{aligned} \quad (21b)$$

Equation (21b) is a *parabolic second-order* PDE, which can be reduced to its canonical form by a proper transformation of variables. In particular, performing the coordinate change

$\tilde{\xi} \triangleq \xi$, $\tilde{s} \triangleq s - \xi$ change, Equation (21b) turns into

$$\frac{\partial^2 u_x(\tilde{\xi}, \tilde{s})}{\partial \tilde{\xi}^2} + \varepsilon_\psi^2 \frac{\varphi^2}{R_r^2} u_x(\tilde{\xi}, \tilde{s}) = -\varepsilon_\psi \frac{\varphi}{R_r} \left[\sigma_y + \frac{\varphi}{R_r} (x_{\mathcal{L}}(\eta) - x_C - \beta) \right], \quad (22)$$

whose solution reads, in terms of ξ and s ,

$$\begin{aligned} u_x(\xi, s) = & -\frac{R_r}{\varepsilon_\psi \varphi} \left[\sigma_y + \frac{\varphi}{R_r} (x_{\mathcal{L}}(\eta) - x_C - \xi) \right] + f(s - \xi) \cos \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi \right) \\ & + g(s - \xi) \sin \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi \right). \end{aligned} \quad (23)$$

Imposing the BC (19) yields the following expression for the displacement vector $\mathbf{u}_r^-(\xi, s)$ for $0 \leq \xi < s \in \mathbb{R}_{\geq 0}$ $0 \leq \xi < s \in \mathbb{R}_{\geq 0}$:

$$\begin{aligned} u_x^-(\xi) = & \frac{\xi}{\varepsilon_\psi} - \left(\frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C}{\varepsilon_\psi} \right) \left[1 - \cos \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi \right) \right] \\ & + \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} \right] \sin \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi \right), \end{aligned} \quad (24a)$$

$$\begin{aligned} u_y^-(\xi) = & \left(\frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C}{\varepsilon_\psi} \right) \sin \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi \right) \\ & + \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} \right] \left[1 - \cos \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi \right) \right]. \end{aligned} \quad (24b)$$

At this point, some elucidation is needed. We said that large values of the steering speed imply a coupling between the slip parameters (σ_x and σ_y) and the spin φ . By looking at Equation (24), this is clear: φ appears in the denominator multiplying both the slip quantities, and also as argument of the trigonometric functions (note that the nonlinear terms are always given by the product of the longitudinal or lateral slip in turn with some function of φ , but never between the translational slips). This effect is not predicted by the classic 1DM, where the bristle displacement is linear in σ_x , σ_y and φ . Furthermore, in the 1DM, the longitudinal component of the deflection depends only on the longitudinal slip and spin, σ_x and φ , whereas the lateral deflection is only a function of σ_y and φ .

The transient solution for (21b) can be computed by performing the second coordinate change $\tilde{\xi} \triangleq \xi - s$, $\tilde{s} \triangleq s$, to obtain

$$\frac{\partial^2 u_x(\tilde{\xi}, \tilde{s})}{\partial \tilde{s}^2} + \varepsilon_\psi^2 \frac{\varphi^2}{R_r^2} u_x(\tilde{\xi}, \tilde{s}) = -\varepsilon_\psi \frac{\varphi}{R_r} \left[\sigma_y + \frac{\varphi}{R_r} (x_{\mathcal{L}}(\eta) - x_C - \alpha - \beta) \right], \quad (25)$$

whose solution reads, in terms of ξ and s ,

$$u_x(\xi, s) = -\frac{R_r}{\varepsilon_\psi \varphi} \left[\sigma_y + \frac{\varphi}{R_r} (x_{\mathcal{L}}(\eta) - x_C - \xi) \right] + f(\xi - s) \cos \left(\varepsilon_\psi \frac{\varphi}{R_r} s \right) + g(\xi - s) \sin \left(\varepsilon_\psi \frac{\varphi}{R_r} s \right). \quad (26)$$

Imposing the IC (20), we get the expressions for the transient deflection $u_t^+(\xi, s)$ of the bristle for $s \leq \xi \leq ls \in \mathbb{R}_{\geq 0} s \leq \xi \leq ls \in \mathbb{R}_{\geq 0}$:

$$u_x^+(\xi, s) = -\frac{R_r}{\varepsilon_\psi \varphi} \sigma_y - \frac{x_{\mathcal{L}}(\eta) - x_C - \xi}{\varepsilon_\psi} + \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - y_C}{\varepsilon_\psi} - u_{y0}(\xi - s) \right] \sin \left(\varepsilon_\psi \frac{\varphi}{R_r} s \right) + \left[u_{x0}(\xi - s) + \frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C - \xi + s}{\varepsilon_\psi} \right] \cos \left(\varepsilon_\psi \frac{\varphi}{R_r} s \right), \quad (27a)$$

$$u_y^+(\xi, s) = \frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - y_C}{\varepsilon_\psi} + \left[u_{x0}(\xi - s) + \frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C - \xi + s}{\varepsilon_\psi} \right] \sin \left(\varepsilon_\psi \frac{\varphi}{R_r} s \right) - \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - y_C}{\varepsilon_\psi} - u_{y0}(\xi - s) \right] \cos \left(\varepsilon_\psi \frac{\varphi}{R_r} s \right). \quad (27b)$$

Similar considerations about the coupling between the slip and spin parameters can be drawn for the transient deflection. In particular, it can be noted that the overall structure of the solution is similar to the one of the steady-state deformation of the bristle, with the difference that the longitudinal coordinate ξ is replaced by the travelled distance s in the trigonometric functions.

3.1. Asymptotic analysis

The previous analysis shows that, for sufficiently large values of ε_ψ and φ , the global solution $u_t(\xi, s)$ is nonlinear in the slip and spin parameters and also every component of the tangential displacement depends on both σ_x and σ_y . Thus, we need to ascertain that Equations (24) and (27) are equivalent to the linear expressions predicted by classic 1DM for small values of ε_ψ and φ . These formulae can be found in many books, although here we refer explicitly to the notation used in [52]. We start from Equation (24) by expanding the

trigonometric functions in Taylor series up to the second term to get

$$u_x^-(\xi) = \frac{\xi}{\varepsilon_\psi} - \left(\frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C}{\varepsilon_\psi} \right) \varepsilon_\psi^2 \frac{\varphi^2}{2R_r^2} \xi^2 + \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} \right] \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi - \varepsilon_\psi^3 \frac{\varphi^3}{6R_r^3} \xi^3 \right), \quad (28a)$$

$$u_y^-(\xi) = \left(\frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C}{\varepsilon_\psi} \right) \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi - \varepsilon_\psi^3 \frac{\varphi^3}{6R_r^3} \xi^3 \right) + \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} \right] \varepsilon_\psi^2 \frac{\varphi^2}{2R_r^2} \xi^2, \quad (28b)$$

and neglecting the linear and higher-order terms in ε_ψ and the quantities x_C, γ_C yields

$$u_x^-(\xi) = \sigma_x \xi - \frac{\varphi}{R_r} \xi \eta, \quad (29a)$$

$$u_y^-(\xi) = \sigma_y \xi + \frac{\varphi}{R_r} \xi \left(x_{\mathcal{L}}(\eta) - \frac{\xi}{2} \right), \quad (29b)$$

which are identical to the formulae derived in [52]. It can be immediately observed that, in (29), both $u_x^-(\xi)$ and $u_y^-(\xi)$ are linear functions of σ_x, σ_y and φ and also $u_x^-(\xi)$ does not depend on the lateral slip and *vice versa*.

Similarly, we perform the series expansion for the transient solution $\mathbf{u}_t^+(\xi, s)$ as follows:

$$u_x^+(\xi, s) = -\frac{R_r}{\varepsilon_\psi \varphi} \sigma_y - \frac{x_{\mathcal{L}}(\eta) - x_C - \xi}{\varepsilon_\psi} + \left[u_{x0}(\xi - s) + \frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C - \xi + s}{\varepsilon_\psi} \right] \left(1 - \varepsilon_\psi^2 \frac{\varphi^2}{2R_r^2} \xi^2 \right) + \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} - u_{y0}(\xi - s) \right] \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi - \varepsilon_\psi^3 \frac{\varphi^3}{6R_r^3} \xi^3 \right), \quad (30a)$$

$$u_y^+(\xi, s) = \frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} + \left[u_{x0}(\xi - s) + \frac{R_r}{\varepsilon_\psi \varphi} \sigma_y + \frac{x_{\mathcal{L}}(\eta) - x_C - \xi + s}{\varepsilon_\psi} \right] \left(\varepsilon_\psi \frac{\varphi}{R_r} \xi - \varepsilon_\psi^3 \frac{\varphi^3}{6R_r^3} \xi^3 \right) - \left[\frac{R_r}{\varepsilon_\psi \varphi} \left(\sigma_x - \frac{1}{\varepsilon_\psi} \right) - \frac{\eta - \gamma_C}{\varepsilon_\psi} - u_{y0}(\xi - s) \right] \left(1 - \varepsilon_\psi^2 \frac{\varphi^2}{2R_r^2} \xi^2 \right), \quad (30b)$$

and neglecting again the terms in ε_ψ and the displacements x_C and y_C we get after some algebra

$$u_x^+(\xi, s) = \sigma_x s - \frac{\varphi}{R_r} \eta s + u_{x0}(\xi - s), \quad (31a)$$

$$u_y^+(\xi, s) = \sigma_y s + \frac{\varphi}{R_r} s \left(x_{\mathcal{L}}(\eta) - \xi + \frac{s}{2} \right) + u_{y0}(\xi - s), \quad (31b)$$

which, again, correspond to the transient solutions found in [52].

It is worth noting that, even though we needed to expand into Taylor series up to the second order, we then neglected the terms in ε_ψ and ε_ψ^2 . This indicates that the classic IDM is very precise only for very small values of ε_ψ , since the highest committed error is in the order of $\mathcal{O}(\varepsilon_\psi)$ (being always $0 \leq \varepsilon_\psi \leq 1$). Of course, this is the case in most vehicle dynamics applications and the two solutions are eventually indistinguishable.

4. Two-dimensional model for large camber angles

The model developed in this section is valid when the camber angle is sufficiently large to generate a non-negligible component of the rolling speed in y direction. However, we consider the case in which the component due to the steering speed is small enough to neglect the coupling between the longitudinal and lateral deflections of the bristle ($\varepsilon_\psi \approx 0$). Also, we assume that the longitudinal component of the velocity field is always negative, i.e. $-\Omega R_r + \Omega \sin \gamma (y - y_C) < 0$, which leads to the following conditions

$$y - y_C - \frac{R_r}{\sin \gamma} < 0 \quad \iff \quad \gamma > 0, \quad \forall y \in \mathcal{P} \quad (32a)$$

$$y - y_C - \frac{R_r}{\sin \gamma} > 0 \quad \iff \quad \gamma < 0, \quad \forall y \in \mathcal{P} \quad (32b)$$

The previous inequalities are fundamental when looking for a particular solution of the problem (see Appendix 1). From a physical perspective, this mathematical requirement states that the tyre semiwidth must be smaller than the lateral coordinate at which the wheel axis intercepts the ground. Of course, in reality this condition is always fulfilled. Another constraint which we impose is that the longitudinal coordinate of the wheel centre must be small if compared with the contact patch length, i.e. $|x_C| < l/2$. Finally, the geometric and slip parameters are again assumed to be constant over the time.

Owing to the previous assumption, Equation (4) can be restated as follows in the adhesion region:

$$\frac{\partial u_x(\mathbf{x}, t)}{\partial t} + \mathbf{v}_t(\mathbf{x}) \cdot \nabla_t u_x(\mathbf{x}, t) = -V_{sx} + \omega_z (y - y_C), \quad \mathbf{x} \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (33a)$$

$$\frac{\partial u_y(\mathbf{x}, t)}{\partial t} + \mathbf{v}_t(\mathbf{x}) \cdot \nabla_t u_y(\mathbf{x}, t) = -V_{sy} - \omega_z (x - x_C), \quad \mathbf{x} \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (33b)$$

where

$$v_x(y) = -\Omega R_r + \Omega \sin \gamma (y - y_C), \quad (34a)$$

$$v_y(x) = -\Omega \sin \gamma (x - x_C), \quad (34b)$$

are the longitudinal and lateral components of the velocity field $\mathbf{v}(\mathbf{x}) = d\mathbf{x}/dt$. In the following, we derive the solutions for the deflection of the bristles for the case of a rectangular contact patch $\mathcal{P} = \{\mathbf{x} \in \mathbb{R}^3 \mid -\frac{l}{2} \leq x \leq \frac{l}{2}, -\frac{w}{2} \leq y \leq \frac{w}{2}, z = 0\}$, where l is the contact patch length and w is the width. Albeit not a very realistic shape for cambered tyres – especially motorcycle tyres, for which large cambering manoeuvres can take place – we consider a rectangular contact shape for two main reasons. The first one relates to the fact that an analytical solution might not be possible for complex geometries; the second one is that the rectangular contact shape is often adopted when dealing with the IDM, and, as previously done with the IDCM, we want to establish a possible connection with the IDM by means of the asymptotic analysis to show a consistency with the results already known in literature.

Recalling the general form of the BC (11), for a rectangular contact patch we can restate them specifically as

$$\text{BCs: } \begin{cases} \mathbf{u}_t \left(\frac{l}{2}, y, t \right) = \mathbf{0}, & -\frac{w}{2} < y < \frac{w}{2}, t \in \mathbb{R}_{>0}, \\ \mathbf{u}_t \left(x, \frac{w}{2} \text{sign}(\gamma), t \right) = \mathbf{0}, & x_C < x < \frac{l}{2}, t \in \mathbb{R}_{>0}, \\ \mathbf{u}_t \left(x, -\frac{w}{2} \text{sign}(\gamma), t \right) = \mathbf{0}, & -\frac{l}{2} < x < x_C, t \in \mathbb{R}_{>0}. \end{cases} \quad (35)$$

Indeed, owing the assumption $\Omega R_r > \Omega \sin \gamma (y - y_C)$, the longitudinal component of the velocity field will be always negative. Hence, all the points located at $x = l/2$ will be inflowing into the contact patch. This can be stated in mathematical terms as

$$\mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x}) \Big|_{x=\frac{l}{2}} < 0 \implies \left\{ \mathbf{x} \in \partial \mathcal{P} \mid x = \frac{l}{2} \right\} \subseteq \mathcal{L}. \quad (36)$$

However, for the case under consideration, we have a two-dimensional velocity field, which means that the bristles can also enter the contact patch from the lateral edges, located at $y = w/2$ and $y = -w/2$, respectively. More specifically, as far as the lateral component $v_y(x)$ is concerned, we can note that, for $\gamma > 0$, $v_y(x) > 0 \iff x < x_C$ and $v_y(x) < 0 \iff x > x_C$. *Vice versa* if the camber angle is negative, i.e. $\gamma < 0$. In formulae:

$$\mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x}) \Big|_{y=\frac{w}{2} \text{sign}(\gamma), x > x_C} < 0 \implies \left\{ \mathbf{x} \in \partial \mathcal{P} \mid y = \frac{w}{2} \text{sign}(\gamma), x > x_C \right\} \subseteq \mathcal{L}, \quad (37)$$

$$\mathbf{v}_t(\mathbf{x}) \cdot \hat{\mathbf{v}}_{\partial \mathcal{P}}(\mathbf{x}) \Big|_{y=-\frac{w}{2} \text{sign}(\gamma), x < x_C} < 0 \implies \left\{ \mathbf{x} \in \partial \mathcal{P} \mid y = -\frac{w}{2} \text{sign}(\gamma), x < x_C \right\} \subseteq \mathcal{L}. \quad (38)$$

In the subsequent analysis, we will show that, because of the three different boundary prescriptions, it is possible to identify three subdomains the contact patch is partitioned in,

which we denote with $\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3$, respectively. These regions correspond to three different analytical expressions for the steady-state deformation of the bristle. Furthermore, as for the classic 1DM and 1DCM, the steady-state and the transient problems can be decoupled since the space boundaries prescribe a constant (actually zero) deformation for the points $\mathbf{x} \in \mathcal{L}$. In particular, the transient solution $\mathbf{u}_t^+(\mathbf{x}, t)$ propagates over the space depending on the specific domain, i.e. on the structure of the speed field. The IC is instead not formally affected by the contact patch shape and reads as in (15).

The detailed derivation of the solution for steady-state conditions is quite cumbersome and is given in Appendix 1. In the next subsections, we limit ourselves to an introductory discussion about its properties and to present the closed-form expressions for the deflections $\mathbf{u}_t^-(\mathbf{x})$ and $\mathbf{u}_t^+(\mathbf{x}, t)$ in each subdomain of \mathcal{P} , with some insights about their physical interpretation.

4.1. Steady-state solution

As for the 1D and 1DC theories, the steady-state and transient problems can be investigated separately. In particular, for the steady-state case, if we assume the existence of a parameter τ , we can write the following system for the independent variables:

$$\frac{dx}{d\tau} = v_x(y) = -\Omega R_r + \Omega \sin \gamma (y - y_C), \quad (39a)$$

$$\frac{dy}{d\tau} = v_y(x) = -\Omega \sin \gamma (x - x_C), \quad (39b)$$

$$\frac{dt}{d\tau} = 1, \quad (39c)$$

whilst for the unknown $u_x(\mathbf{x}, t)$ and $u_y(\mathbf{x}, t)$ we have

$$\frac{du_x(\mathbf{x}, t)}{d\tau} = -V_{sx} + \omega_z (y - y_C), \quad (40a)$$

$$\frac{du_y(\mathbf{x}, t)}{d\tau} = -V_{sy} - \omega_z (x - x_C). \quad (40b)$$

Combining (39a) with (39b) and (39a) with (39c), respectively, provides the characteristics lines for the independent variables:

$$c_1 = (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2, \quad (41a)$$

$$c_2 = t - \frac{1}{\Omega \sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right). \quad (41b)$$

Two important considerations can be drawn by looking at the previous relations. From Equation (41a), we deduce that a first set of characteristic lines is given by a family of circumferences whose centre is located at $C_\gamma \triangleq (x_C, y_C + \frac{R_r}{\sin \gamma})$, which is exactly the point at which the wheel axis intercepts the *Oxy* plane. We call this point *cambering centre*. The radius at $\mathbf{x} = \mathbf{0}$ reads $R_\gamma = \sqrt{x_C^2 + (y_C + \frac{R_r}{\sin \gamma})^2}$. The second Equation (41b) provides an

explicit relation for the time as a function of the space coordinates; it is worth to emphasise that (41a) does not involve the time, which implies that the steady-state solution can be sought independently of the transient one (basically, (41b) is redundant when looking for the steady-state solution).

The general solution for the vector displacement $\mathbf{u}_t^-(\mathbf{x})$ can be then found by combining (39a) with (40a) and (40b), separately, to obtain

$$c_{3x} = u_x^-(\mathbf{x}, t) - \sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) + \frac{\varphi}{\sin \gamma} \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) + x \right], \quad (42a)$$

$$c_{3y} = u_y^-(\mathbf{x}, t) - \sigma_y \frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) - \frac{\varphi}{\sin \gamma} \left(\frac{R_r}{\sin \gamma} + y_C - y \right). \quad (42b)$$

Invoking the IFT, Equation (41) with (42a) and (42b), respectively, admit a solution in the forms

$$u_x^-(\mathbf{x}, t) = \sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) - \frac{\varphi}{\sin \gamma} \times \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) + x \right] + G_x(c_1(\mathbf{x}), c_2(\mathbf{x}, t)), \quad (43a)$$

$$u_y^-(\mathbf{x}, t) = \sigma_y \frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) + \frac{\varphi}{\sin \gamma} \left(\frac{R_r}{\sin \gamma} + y_C - y \right) + G_y(c_1(\mathbf{x}), c_2(\mathbf{x}, t)), \quad (43b)$$

where $G_x(\cdot, \cdot)$ and $G_y(\cdot, \cdot)$ are arbitrary functions whose arguments read as in (41).

Now, the particular solution will depend on the BCs. At this point, an important reflection concerning the number of different solutions which can be obtained is needed. As we mentioned previously, (33) are two *transport equations* which only require one BC to be completely determined. Since from (35) we have three different BCs, we need to impose each of them separately and look for three independent solutions which apply to three areas \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 of the contact patch. These regions are given by the following domains (see Appendix 2):

$$\mathcal{P}_1 \triangleq \{ \mathcal{P} \setminus (\mathcal{P}_2 \cup \mathcal{P}_3) \}, \quad (44a)$$

$$\mathcal{P}_2 \triangleq \{ \mathbf{x} \in \mathcal{P} \mid R_1^2 < \Gamma(\mathbf{x}) < R_2^2 \}, \quad (44b)$$

$$\mathcal{P}_3 \triangleq \{ \mathbf{x} \in \mathcal{P} \mid R_3^2 < \Gamma(\mathbf{x}) < R_4^2 \cap x < x_C \}, \quad (44c)$$

where $\Gamma(\mathbf{x})$ reads

$$\Gamma(\mathbf{x}) \triangleq c_1(\mathbf{x}) = (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2, \quad (45)$$

and the radii are as follows:

$$R_0 \triangleq \frac{l}{2} - x_C, \quad (46a)$$

$$R_1 \triangleq \frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma}, \quad (46b)$$

$$R_2 \triangleq \sqrt{R_1^2 + R_0^2}, \quad (46c)$$

$$R_3 \triangleq \frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}, \quad (46d)$$

$$R_4 \triangleq \sqrt{R_3^2 + R_5^2}, \quad (46e)$$

$$R_5 \triangleq \frac{l}{2} + x_C. \quad (46f)$$

Note that, owing the assumptions (32) and $|x_C| < l/2$, the radii R_1 and R_3 can be negative, whilst R_0 and R_5 are always positive.

Figure 2 shows the three domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 for different values of the camber angle $\gamma > 0$ and for $x_C = y_C = 0$. More specifically, the main green area represents the region \mathcal{P}_1 , whilst the red and blue ones are \mathcal{P}_2 and \mathcal{P}_3 , respectively. It can be seen that larger values of the camber angles cause the two subdomains \mathcal{P}_2 and \mathcal{P}_3 to be more extended. Of course, for $\gamma = 0$, the theory reduces to the standard one (see Section 4.3). The black arrows represent the velocity field at each point of the contact patch and are always tangent to circles of arbitrary radius whose centre is located in C_γ . In particular, it can be seen that the arrows located at the right edge of the contact patch ($x = l/2$) point inside, since the bristles are inflowing into the contact area; Analogously, other bristles enter the contact patch from the two lateral edges for $x > x_C$ and $x < x_C$, according to the corresponding BCs. Depending on the specific BC which applies in turn, the steady-state tangential deformation of the bristle $\mathbf{u}_t^-(\mathbf{x})$ will be given by a different expression.

Hence, we denote with $\mathbf{u}_{1t}^-(\mathbf{x})$, $\mathbf{u}_{2t}^-(\mathbf{x})$, $\mathbf{u}_{3t}^-(\mathbf{x})$ the steady-state solutions in the three different domains. We have

$$u_{1x}^-(\mathbf{x}) = \sigma_x \Sigma_1(\mathbf{x}) + \varphi \Phi_{1x}(\mathbf{x}), \quad \mathbf{x} \in \mathcal{P}_1^-, s \in \mathbb{R}_{\geq 0}, \quad (47a)$$

$$u_{2x}^-(\mathbf{x}) = \sigma_x \Sigma_2(\mathbf{x}) + \varphi \Phi_{2x}(\mathbf{x}), \quad \mathbf{x} \in \mathcal{P}_2^-, s \in \mathbb{R}_{\geq 0}, \quad (47b)$$

$$u_{3x}^-(\mathbf{x}) = \sigma_x \Sigma_3(\mathbf{x}) + \varphi \Phi_{3x}(\mathbf{x}), \quad \mathbf{x} \in \mathcal{P}_3^-, s \in \mathbb{R}_{\geq 0}, \quad (47c)$$

and

$$u_{1y}^-(\mathbf{x}) = \sigma_y \Sigma_1(\mathbf{x}) + \varphi \Phi_{1y}(\mathbf{x}), \quad \mathbf{x} \in \mathcal{P}_1^-, s \in \mathbb{R}_{\geq 0}, \quad (48a)$$

$$u_{2y}^-(\mathbf{x}) = \sigma_y \Sigma_2(\mathbf{x}) + \varphi \Phi_{2y}(\mathbf{x}), \quad \mathbf{x} \in \mathcal{P}_2^-, s \in \mathbb{R}_{\geq 0}, \quad (48b)$$

$$u_{3y}^-(\mathbf{x}) = \sigma_y \Sigma_3(\mathbf{x}) + \varphi \Phi_{3y}(\mathbf{x}), \quad \mathbf{x} \in \mathcal{P}_3^-, s \in \mathbb{R}_{\geq 0}, \quad (48c)$$

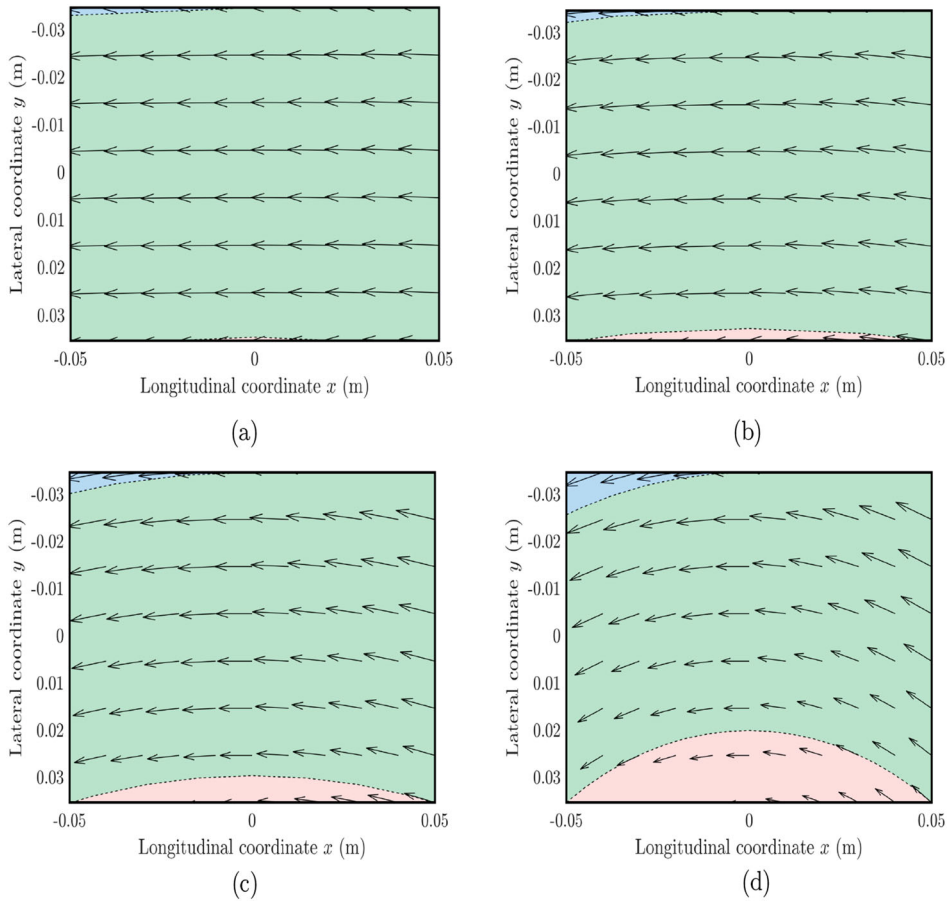


Figure 2. Subdomains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 of the contact patch \mathcal{P} (represented by the rectangles) for different values of the camber angle $\gamma = 10, 30, 50$ and 70° , respectively. The larger green area corresponds to the domain \mathcal{P}_1 , whilst the red and blue one to \mathcal{P}_2 and \mathcal{P}_3 . The arrows represent the velocity field $\mathbf{v}_t(\mathbf{x})$ and are always tangent to the trajectories of the bristles, which coincide with the characteristics lines given by (41a). The length and the width of the contact patch are $l = 0.1$ and $w = 0.07$ (m). (a) Domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 for a camber angle of $\gamma = 10^\circ$. (b) Domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 for a camber angle of $\gamma = 30^\circ$. (c) Domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 for a camber angle of $\gamma = 50^\circ$. (d) Domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 for a camber angle of $\gamma = 70^\circ$.

where $\Sigma(\mathbf{x})$ and $\Phi(\mathbf{x})$ are as follows:

$$\Sigma_1(\mathbf{x}) \triangleq \frac{R_r}{\sin \gamma} \left[\Xi(\mathbf{x}) + \arctan \left(\frac{R_0}{\sqrt{\Gamma(\mathbf{x}) - R_0^2}} \right) \text{sign}(\gamma) \right], \quad (49a)$$

$$\Sigma_2(\mathbf{x}) \triangleq \frac{R_r}{\sin \gamma} \left[\Xi(\mathbf{x}) - \arctan \left(\frac{\sqrt{\Gamma(\mathbf{x}) - R_1^2}}{R_1} \right) \right], \quad (49b)$$

$$\Sigma_3(\mathbf{x}) \triangleq \frac{R_r}{\sin \gamma} \left[\Xi(\mathbf{x}) - \arctan \left(\frac{\sqrt{\Gamma(\mathbf{x}) - R_3^2}}{R_3} \right) \right], \quad (49c)$$

$$\Phi_{1x}(\mathbf{x}) \triangleq -\frac{1}{\sin \gamma} \left(x - \frac{l}{2} + \Sigma_1(\mathbf{x}) \right), \quad (49d)$$

$$\Phi_{2x}(\mathbf{x}) \triangleq -\frac{1}{\sin \gamma} \left(x - x_C - \sqrt{\Gamma(\mathbf{x}) - R_1^2} + \Sigma_2(\mathbf{x}) \right), \quad (49e)$$

$$\Phi_{3x}(\mathbf{x}) \triangleq -\frac{1}{\sin \gamma} \left(x - x_C + \sqrt{\Gamma(\mathbf{x}) - R_3^2} + \Sigma_3(\mathbf{x}) \right), \quad (49f)$$

$$\Phi_{1y}(\mathbf{x}) \triangleq \frac{1}{\sin \gamma} \left(\frac{R_r}{\sin \gamma} + y_C - y - \sqrt{\Gamma(\mathbf{x}) - R_0^2} \text{sign}(\gamma) \right), \quad (49g)$$

$$\Phi_{2y}(\mathbf{x}) \triangleq \frac{1}{\sin \gamma} \left(\frac{w}{2} \text{sign}(\gamma) - y \right), \quad (49h)$$

$$\Phi_{3y}(\mathbf{x}) \triangleq -\frac{1}{\sin \gamma} \left(\frac{w}{2} \text{sign}(\gamma) + y \right), \quad (49i)$$

in which we have defined

$$\Xi(\mathbf{x}) \triangleq \arctan \left(\frac{x - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right). \quad (50)$$

An interesting result is that the two solutions $\mathbf{u}_{1t}^-(\mathbf{x})$ and $\mathbf{u}_{2t}^-(\mathbf{x})$ are continuous at the interface between \mathcal{P}_1 and \mathcal{P}_2 . Indeed, it is easy to verify that the following relations are satisfied:

$$\Sigma_1(\mathbf{x}) \Big|_{\mathcal{C}_2(\mathbf{x})} = \Sigma_2(\mathbf{x}) \Big|_{\mathcal{C}_2(\mathbf{x})}, \quad (51a)$$

$$\Phi_{1x}(\mathbf{x}) \Big|_{\mathcal{C}_2(\mathbf{x})} = \Phi_{2x}(\mathbf{x}) \Big|_{\mathcal{C}_2(\mathbf{x})}, \quad (51b)$$

$$\Phi_{1y}(\mathbf{x}) \Big|_{\mathcal{C}_2(\mathbf{x})} = \Phi_{2y}(\mathbf{x}) \Big|_{\mathcal{C}_2(\mathbf{x})}. \quad (51c)$$

From a mathematical viewpoint, this structure of this solution is a consequence of the fact $\mathbf{u}_{1t}^-(\frac{l}{2}, \frac{w}{2}, t) = \mathbf{u}_{2t}^-(\frac{l}{2}, \frac{w}{2}, t) = 0$. Conversely, a discontinuity occurs between the regions \mathcal{P}_1 and \mathcal{P}_3 . This can be explained by looking at the physics underlying the rolling of the tyre. On the one hand, the bristles inflowing in \mathcal{P}_3 are the ones which had outflowed from the region \mathcal{P}_1 and enter again undeformed at the edge $-w/2$; on the other hand, the bristles of \mathcal{P}_1 which are not transported outside the contact patch keep rolling and their deformation rapidly increases over the travelled distance. We point out, however, that in reality it is unlikely that the contact patch conserves its rectangular shape for large camber angles, whilst there is some evidence in literature that curved geometries can be observed

[33,34,61]. We may then conjecture that the third region \mathcal{P}_3 disappears; in this case, the deformation would always be a continuous function in space.

4.2. Transient solution

To seek for the transient solution, it is more convenient to express all the dependent variables as a function of the time. In order to do that, we derive Equation (39b) with respect to the parameter τ , or, equivalently, to the time, and then substitute Equation (39a), to get

$$x(t) = (x_0 - x_C) \cos(\Omega \sin \gamma t) + \left(y_0 - y_C - \frac{R_r}{\sin \gamma} \right) \sin(\Omega \sin \gamma t) + x_C, \quad (52a)$$

$$y(t) = -(x_0 - x_C) \sin(\Omega \sin \gamma t) + \left(y_0 - y_C - \frac{R_r}{\sin \gamma} \right) \cos(\Omega \sin \gamma t) + y_C + \frac{R_r}{\sin \gamma}. \quad (52b)$$

Then, integrating (40a) and (40b) with respect to t and renaming $s \triangleq \Omega R_r t$ yields

$$u_x^+(\mathbf{x}, s) = \sigma_x s - \frac{\varphi}{\sin \gamma} (x - x_0(\mathbf{x}, s) + s) + u_{x0}(\mathbf{x}_0(\mathbf{x}, s)), \quad \mathbf{x} \in \mathcal{P}^+, \quad (53a)$$

$$u_y^+(\mathbf{x}, s) = \sigma_y s - \frac{\varphi}{\sin \gamma} (y - y_0(\mathbf{x}, s) + s) + u_{y0}(\mathbf{x}_0(\mathbf{x}, s)), \quad \mathbf{x} \in \mathcal{P}^+, \quad (53b)$$

where the initial datum (ID) $\mathbf{x}_0(s) = [x_0(s) \ y_0(s)]^T$ can be obtained by inverting (52)

$$x_0(\mathbf{x}, s) = (x - x_C) \cos\left(\frac{\sin \gamma}{R_r} s\right) - \left(y - y_C - \frac{R_r}{\sin \gamma}\right) \sin\left(\frac{\sin \gamma}{R_r} s\right) + x_C, \quad (54a)$$

$$y_0(\mathbf{x}, s) = (x - x_C) \sin\left(\frac{\sin \gamma}{R_r} s\right) + \left(y - y_C - \frac{R_r}{\sin \gamma}\right) \cos\left(\frac{\sin \gamma}{R_r} s\right) + y_C + \frac{R_r}{\sin \gamma}. \quad (54b)$$

We note that, in this case, the notion of *travelled distance* does not apply directly to the quantity s . Indeed, ΩR_r only represents the main component of the angular speed. However, if $x_C = y_C = 0$, s may be still interpreted as an averaged *travelled distance* (note that ΩR_r is exactly the mean value of the speed field over the contact patch).

Equations (53a) and (53b) represent the general transient solution for the planar vector displacement $\mathbf{u}_t^+(\mathbf{x}, s)$, whereas a particular solution can be determined by imposing the IC, i.e. an explicit expression for $\mathbf{u}_t(\mathbf{x}, 0) = \mathbf{u}_{t0}(\mathbf{x})$. Generally speaking, this IC does not depend on the specific subdomain of \mathcal{P} and thus the transient solution is a global one. However, because of the discontinuity in the steady-state solution between the domains \mathcal{P}_1 and \mathcal{P}_3 , the IC must also be applied separately, i.e. we need $\mathbf{u}_{1t0}(\mathbf{x})$, $\mathbf{u}_{2t0}(\mathbf{x})$ and $\mathbf{u}_{3t0}(\mathbf{x})$ in \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 , respectively. We also need to identify the portion of the space which the transient solution applies to. This investigation can be conducted with respect to each subdomain of \mathcal{P} . In order to do this, we can look at (39c) and seek for a solution depending on the BC. In particular, from (41b), we deduce that the following conditions must be respectively

satisfied for the regions \mathcal{P}_1^+ , \mathcal{P}_2^+ and \mathcal{P}_3^+ :

$$\Sigma_1(\mathbf{x}) \geq s, \quad (55a)$$

$$\Sigma_2(\mathbf{x}) \geq s, \quad (55b)$$

$$\Sigma_3(\mathbf{x}) \geq s. \quad (55c)$$

Thus, we can define the following subdomains:

$$\mathcal{P}_1^- \triangleq \{\mathbf{x} \in \mathcal{P}_1 \mid \Sigma_1(\mathbf{x}) < s\}, \quad (56a)$$

$$\mathcal{P}_2^- \triangleq \{\mathbf{x} \in \mathcal{P}_2 \mid \Sigma_2(\mathbf{x}) < s\}, \quad (56b)$$

$$\mathcal{P}_3^- \triangleq \{\mathbf{x} \in \mathcal{P}_3 \mid \Sigma_3(\mathbf{x}) < s\}, \quad (56c)$$

$$\mathcal{P}_1^+ \triangleq \{\mathbf{x} \in \mathcal{P}_1 \mid \Sigma_1(\mathbf{x}) \geq s\}, \quad (56d)$$

$$\mathcal{P}_2^+ \triangleq \{\mathbf{x} \in \mathcal{P}_2 \mid \Sigma_2(\mathbf{x}) \geq s\}, \quad (56e)$$

$$\mathcal{P}_3^+ \triangleq \{\mathbf{x} \in \mathcal{P}_3 \mid \Sigma_3(\mathbf{x}) \geq s\}, \quad (56f)$$

so that the general solution for the planar vector displacements reads

$$\mathbf{u}_t(\mathbf{x}, s) = \begin{cases} \mathbf{u}_{1t}^-(\mathbf{x}), & \mathbf{x} \in \mathcal{P}_1^-, s \in \mathbb{R}_{\geq 0}, \\ \mathbf{u}_{2t}^-(\mathbf{x}), & \mathbf{x} \in \mathcal{P}_2^-, s \in \mathbb{R}_{\geq 0}, \\ \mathbf{u}_{3t}^-(\mathbf{x}), & \mathbf{x} \in \mathcal{P}_3^-, s \in \mathbb{R}_{\geq 0}, \\ \mathbf{u}_{1t}^+(\mathbf{x}, s), & \mathbf{x} \in \mathcal{P}_1^+, s \in \mathbb{R}_{\geq 0}, \\ \mathbf{u}_{2t}^+(\mathbf{x}, s), & \mathbf{x} \in \mathcal{P}_2^+, s \in \mathbb{R}_{\geq 0}, \\ \mathbf{u}_{3t}^+(\mathbf{x}, s), & \mathbf{x} \in \mathcal{P}_3^+, s \in \mathbb{R}_{\geq 0}. \end{cases} \quad (57)$$

Finally, it is easy to verify that, in each subdomain, the following relations hold for $s = \Sigma(\mathbf{x})$. Indeed, it is

$$\mathbf{u}_{1t}^+(\mathbf{x}, \Sigma_1(\mathbf{x})) = \mathbf{u}_{1t}^-(\mathbf{x}), \quad (58a)$$

$$\mathbf{u}_{2t}^+(\mathbf{x}, \Sigma_2(\mathbf{x})) = \mathbf{u}_{2t}^-(\mathbf{x}), \quad (58b)$$

$$\mathbf{u}_{3t}^+(\mathbf{x}, \Sigma_3(\mathbf{x})) = \mathbf{u}_{3t}^-(\mathbf{x}), \quad (58c)$$

since, after some manipulation, we can find that $x_0(\Sigma_1(\mathbf{x})) = \frac{l}{2}$, $y_0(\Sigma_2(\mathbf{x})) = \frac{w}{2} \text{sign}(\gamma)$ and $y_0(\Sigma_3(\mathbf{x})) = -\frac{w}{2} \text{sign}(\gamma)$, which implies that the displacements $\mathbf{u}_{t0}(\mathbf{x})$ vanish when evaluated in $\Sigma(\mathbf{x})$ because of the boundary prescription. Roughly speaking, this means that the solutions are always continuous at the interface between the steady-state domains and the corresponding transient ones. Furthermore, from the first three of (51), we deduce that the curve that makes the transition between the two regimes is continuous at the interface between the subdomains \mathcal{P}_1 and \mathcal{P}_2 .

Figure 3 shows the time evolution of the steady-state and transient domains in the contact patch for different values of the nondimensional travelled distance \bar{s} (normalised with respect to the total length) and $x_C = y_C = 0$. It can be seen that the separation lines which represent the interface between the steady-state and transient solutions of \mathcal{P}_1 and \mathcal{P}_2

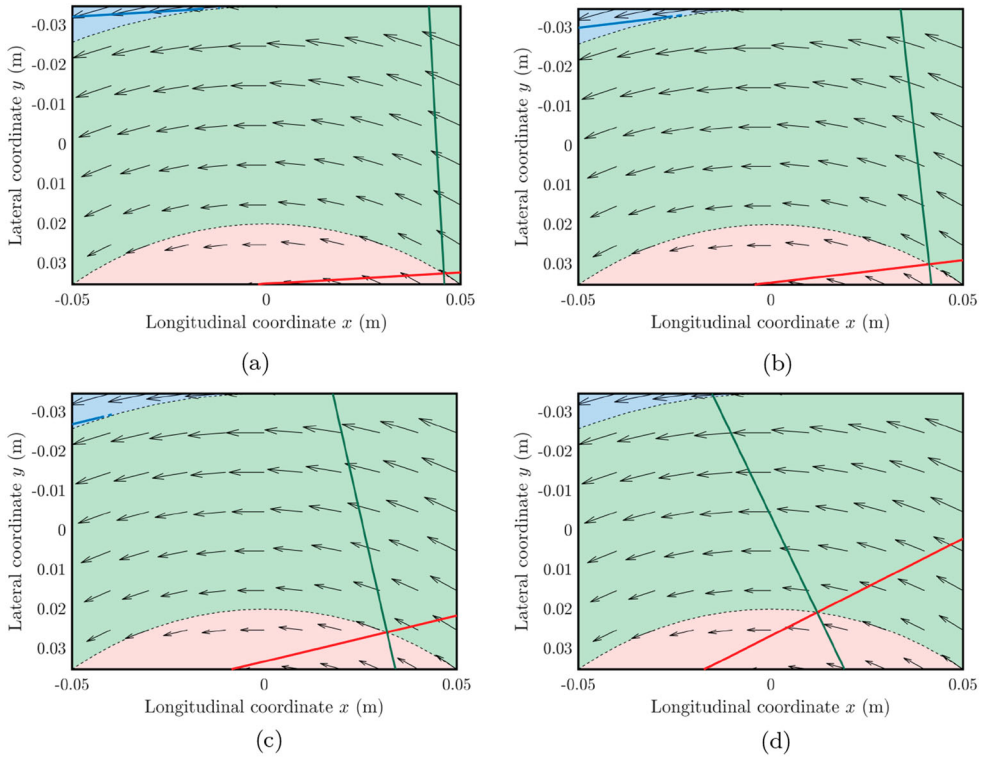


Figure 3. Steady-state and transient domains for each region of \mathcal{P} for a constant value of the camber angle $\gamma = 70^\circ$ and different travelled distances $\bar{s} = 1/16, 1/8, 1/4$ and $1/2$. The green, red and blue solid lines represent the interface between the steady-state and transient solution in the domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 , respectively. More specifically, the steady-state regions correspond to the right portion of the green and red areas for \mathcal{P}_1 and \mathcal{P}_2 , and to the upper part of the blue domain for \mathcal{P}_3 . It can be noted that steady-state conditions are reached in the region \mathcal{P}_3 relatively faster (the transient extinguishes in the blue area already for $\bar{s} = 1/2$). The length and the width of the contact patch are $l = 0.1$ and $w = 0.07$ (m). (a) Steady-state and transient domains for each region of \mathcal{P} for $\bar{s} = 1/16$. (b) Steady-state and transient domains for each region of \mathcal{P} for $\bar{s} = 1/8$. (c) Steady-state and transient domains for each region of \mathcal{P} for $\bar{s} = 1/4$. (d) Steady-state and transient domains for each region of \mathcal{P} for $\bar{s} = 1/2$.

(green and red dashed curves, respectively) are continuous on the circumference $\mathcal{C}_2(\mathbf{x})$ (see Appendix 2). For the region \mathcal{P}_1 , it can be also noted that the bristles located in the negative side of the contact patch reach the steady-state conditions faster. This is due to the fact that, for $y < 0$, the main component of the rolling speed and the variable part have are concordant. Conversely, for positive values of y the total longitudinal speed decreases in absolute value because its given by two opposite contributions.

On the other hand, in the domain \mathcal{P}_2 the bristle inflowing in the contact patch (located at the edge $y = \frac{w}{2}\text{sign}(\gamma)$) are always in steady-state conditions and the transient extinguishes faster in the lower part of the region. This can be explained by considering that the solution is basically given by a wave travelling in space: as the first bristles enter the contact patch, they are propagated towards the trailing edge and steady-state conditions immediately take place, despite the lower values of the speed.

With the same reasoning it is possible to explain what happens in the domain \mathcal{P}_3 . In this case, the bristles are entering the contact patch from the edge $y = -\frac{w}{2}\text{sign}(\gamma)$, and hence the steady-state solution propagates from the upper to the lower part of the blue region. Furthermore, in this case the transient extinguishes even faster, since the rolling speed in the upper side is characterised by a larger longitudinal component in absolute value.

An alternative geometrical interpretation is finally depicted in Figure 4, where $\Sigma_1(\mathbf{x})$, $\Sigma_2(\mathbf{x})$ and $\Sigma_3(\mathbf{x})$ are plotted for two different values of the nondimensional travelled distance $\bar{s} = 1/8$ and $1/2$, respectively. It can be noted that the separation lines between the steady-state and transient subdomains in each region of the contact patch are tangent to the circumferences $\mathcal{C}_0(\mathbf{x})$, $\mathcal{C}_1(\mathbf{x})$ and $\mathcal{C}_3(\mathbf{x})$ (see Appendix 2 for a more detailed discussion). These circumferences represent the domains for the quantities under the square roots in (49a), (49b) and (49c).

4.3. Asymptotic analysis

As for the nonlinear case, we need to ascertain that the two-dimensional theory is consistent with the results predicted by the classic one for small values of the camber angle γ and $x_C = y_C = 0$. Starting from the steady-state solution, it is immediate to verify that the two subdomains \mathcal{P}_2 and \mathcal{P}_3 vanish for γ sufficiently small. Indeed, from Equations (44b) and (44c), we get $\mathcal{P}_2 = \mathcal{P}_3 \approx \emptyset$ as in the standard brush models. Thus, we only have to analyse the solution $\mathbf{u}_{1f}^-(\mathbf{x})$.

In particular, by expanding into Taylor series the square root and the trigonometric functions in (49a) and neglecting the term y , we get

$$\Sigma_1(\mathbf{x}) \approx \frac{l}{2} - x. \quad (59)$$

The above result also makes allowance for some consideration about the two different areas of \mathcal{P} which the steady-state and the transient solution apply to. Indeed, since we have defined \mathcal{P}_1^- and \mathcal{P}_1^+ , Equation (59) states that, if the camber angle γ is sufficiently small, the two subdomains are separated by the straight line $x = l/2 - s$ and the *travelled distance* regains its original meaning.

For $\Phi_{1x}(\mathbf{x})$, we expand the trigonometric functions up to the second term and then approximate

$$\frac{1}{y - \frac{R_r}{\sin \gamma}} \approx -\frac{\sin \gamma}{R_r} - \frac{\sin^2 \gamma y}{R_r^2}$$

to get

$$\Phi_{1x}(\mathbf{x}) \approx -\frac{1}{R_r} \left(\frac{l}{2} - x \right) y. \quad (60)$$

Finally, for $\Phi_{1y}(\mathbf{x})$ we perform the Taylor expansion of the square root up to the second order yielding

$$\Phi_{1y}(\mathbf{x}) \approx \frac{1}{2R_r} \left(\frac{l^2}{4} - x^2 \right). \quad (61)$$

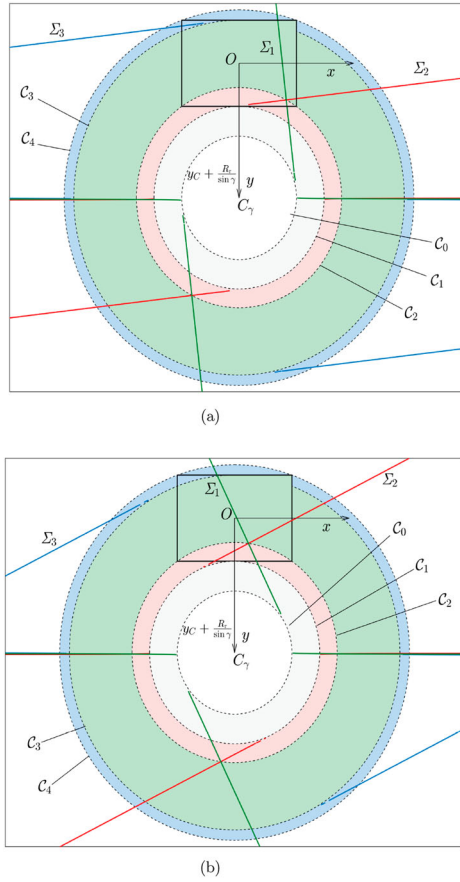


Figure 4. Steady-state and transient domains for each region of \mathcal{P} for a constant value of the camber angle $\gamma = 70^\circ$ and different travelled distances $\bar{s} = 1/8$ and $1/2$. The functions Σ_1 , Σ_2 and Σ_3 correspond to the green, red and blue curve, respectively, and represent the interface between the steady-state and blacktransient blacksolution in the domains \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 . It can be seen that they rotate tangentially to the circumferences C_0 , C_1 and C_3 . (a) Steady-state and transient subdomains of the contact patch \mathcal{P} corresponding to a normalised travelled distance of $\bar{s} = 1/8$. (b) Steady-state and transient subdomains of the contact patch \mathcal{P} corresponding to a normalised travelled distance of $\bar{s} = 1/2$.

Note that, with $\Sigma_1(\mathbf{x})$, $\Phi_{1x}(\mathbf{x})$ and $\Phi_{1y}(\mathbf{x})$ defined respectively as in (59), (60) and (61) and performing the change of variables $\xi = l/2 - x$ and $y = \eta$, Equations (47a) and (48a) reduce exactly to (29a).

Now we extend the analysis to the transient solution. We have to find an approximate expression for the ID $\mathbf{x}_0(\mathbf{x}, s)$. Starting from $x_0(\mathbf{x}, s)$, we obtain, for small values of γ ,

$$x_0(\mathbf{x}, s) \approx x - y \frac{\sin \gamma}{R_r} s + s. \tag{62}$$

Analogously, Taylor expansion for $y_0(s)$ provides

$$y_0(\mathbf{x}, s) \approx y + \frac{\sin \gamma}{R_r} s \left(x + \frac{s}{2} \right). \tag{63}$$

Combining (53a) and (53b) with the previous approximate expressions (62) and (63) with the usual change of variables $\xi = l/2 - x$ and $y = \eta$ gives again (31a) and (31b).

5. General theory

Finally, we consider the complete model which accounts for both the coupling between the slips and the spin parameter and for the presence of a two-dimensional velocity field in the contact patch. Since this model combines both the previous ones, we can already infer some properties of the solution. In the first instance, for a rectangular contact patch, we can expect to find the same three regions of the contact patch, namely \mathcal{P}_1 , \mathcal{P}_2 , \mathcal{P}_3 , in which the solution for the deflection of the bristle will be given by a different analytical expression. This analogy with the 2DM is a consequence of the fact that the BCs are only determined by the structure of the vector field (or, equivalently, the characteristic lines of the problem), which is not influenced by the *steering ratio* ε_ψ . On the other hand, we may anticipate that each tangential component of the deformation $\mathbf{u}_t(\mathbf{x}, t)$ will depend on both the translational slips σ_x and σ_y , and that slip parameters will multiply nonlinear functions of the spin φ , which are likely to be elementary trigonometric functions.

We start by introducing again the tyre-road contact equations in the following form:

$$\frac{\partial u_x(\mathbf{x}, t)}{\partial t} = D_t[u_x(\mathbf{x}, t)] - V_{sx} + \omega_z(y - y_C + \varepsilon_\psi u_y(\mathbf{x}, t)), \quad \mathbf{x} \in \overset{\circ}{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (64a)$$

$$\frac{\partial u_y(\mathbf{x}, t)}{\partial t} = D_t[u_y(\mathbf{x}, t)] - V_{sy} - \omega_z(x - x_C + \varepsilon_\psi u_x(\mathbf{x}, t)), \quad \mathbf{x} \in \overset{\circ}{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (64b)$$

in which the tangential partial differential operator D_t has been defined reading

$$D_t \triangleq -\mathbf{v}_t(\mathbf{x}) \cdot \nabla_t = -\left(v_x(y) \frac{\partial}{\partial x} + v_y(x) \frac{\partial}{\partial y} \right), \quad (65)$$

where the speeds $v_x(y)$ and $v_y(x)$ are exactly as in (34).

Similarly to [62], we can now look for a general solution in the form

$$u_x(\mathbf{x}, t) = \phi_1(t)U_x(\mathbf{x}, t) + \phi_2(t)U_y(\mathbf{x}, t) + \tilde{u}_x(x), \quad (66a)$$

$$u_y(\mathbf{x}, t) = \psi_1(t)U_x(\mathbf{x}, t) + \psi_2(t)U_y(\mathbf{x}, t) + \tilde{u}_y(y), \quad (66b)$$

where $\tilde{u}_x(x)$ and $\tilde{u}_y(y)$ can be cleverly assigned to reduce Equation (64) to homogeneous form. In particular, the choice

$$\tilde{u}_x(x) = -x + x_C - \frac{V_{sy}}{\varepsilon_\psi \omega_z}, \quad (67a)$$

$$\tilde{u}_y(y) = -y + y_C + \frac{V_{sx} + \Omega R_r}{\varepsilon_\psi \omega_z}, \quad (67b)$$

or, equivalently, in terms of the slip parameters

$$\tilde{u}_x(x) = -x + x_C - \frac{\sigma_y R_r}{\varepsilon_\psi \varphi}, \quad (68a)$$

$$\tilde{u}_y(y) = -y + y_C + \frac{\sigma_x R_r}{\varepsilon_\psi \varphi} - \frac{R_r}{\varepsilon_\psi \varphi}, \quad (68b)$$

turns (64) into the following set of PDEs and ODEs:

$$\frac{\partial U_x(\mathbf{x}, t)}{\partial t} + v_x(y) \frac{\partial U_x(\mathbf{x}, t)}{\partial x} + v_y(x) \frac{\partial U_x(\mathbf{x}, t)}{\partial y} = 0, \quad \mathbf{x} \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (69a)$$

$$\frac{\partial U_y(\mathbf{x}, t)}{\partial t} + v_x(y) \frac{\partial U_y(\mathbf{x}, t)}{\partial x} + v_y(x) \frac{\partial U_y(\mathbf{x}, t)}{\partial y} = 0, \quad \mathbf{x} \in \mathring{\mathcal{P}}, \quad t \in \mathbb{R}_{>0}, \quad (69b)$$

together with

$$\dot{\phi}(t) = \varepsilon_\psi \omega_z \psi(t), \quad t \in \mathbb{R}_{>0}, \quad (70a)$$

$$\dot{\psi}(t) = -\varepsilon_\psi \omega_z \phi(t), \quad t \in \mathbb{R}_{>0}, \quad (70b)$$

where we have introduced Newton's notation for the time derivative. Note that now $\phi_1(t)$, $\psi_1(t)$ and $\phi_2(t)$, $\psi_2(t)$ in Equation (66) are two independent solutions of (70a) and (70b).

Now we draw some considerations about the general structure of the solution. Let us start from analysing the first part of the right term of Equation (66). From (70a) and (70b), we find

$$\phi(t) = a_1 \cos(\varepsilon_\psi \omega_z t) + a_2 \sin(\varepsilon_\psi \omega_z t), \quad (71a)$$

$$\psi(t) = -a_1 \sin(\varepsilon_\psi \omega_z t) + a_2 \cos(\varepsilon_\psi \omega_z t). \quad (71b)$$

Furthermore, by means of the method of the characteristics, we are able to solve (69a) and (69b), which provide

$$U_x(\mathbf{x}, t) = G_x(c_1(\mathbf{x}), c_2(\mathbf{x}, t)), \quad (72a)$$

$$U_y(\mathbf{x}, t) = G_y(c_1(\mathbf{x}), c_2(\mathbf{x}, t)), \quad (72b)$$

since the characteristics curve are exactly the same as for the linear two-dimensional case. Thus, we can restate (66) as

$$u_x(\mathbf{x}, t) = A_x(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) \cos(\varepsilon_\psi \omega_z t) + A_y(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) \sin(\varepsilon_\psi \omega_z t) + \tilde{u}_x(x), \quad (73a)$$

$$u_y(\mathbf{x}, t) = -A_x(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) \sin(\varepsilon_\psi \omega_z t) + A_y(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) \cos(\varepsilon_\psi \omega_z t) + \tilde{u}_y(y), \quad (73b)$$

where we have denoted with $A_x(\cdot, \cdot) \triangleq a_1 G_x(\cdot, \cdot)$ and $A_y(\cdot, \cdot) \triangleq a_2 G_y(\cdot, \cdot)$ ⁷. The steady-state solution $\mathbf{u}_t^-(\mathbf{x}, t)$ can be obtained again by imposing the BCs. In particular, we assume

that, with some caution on the invertibility, it is possible to find a parametrisation $f(\cdot)$ ⁸ for the points $\mathbf{x} \in \mathcal{L}$ so that

$$x = f(y), \quad (74a)$$

$$\begin{aligned} c_1(\mathbf{x}) &= (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2 \\ &= (f(y) - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2 \implies y = g(c_1), \end{aligned} \quad (74b)$$

$$c_2(\mathbf{x}, t) = t - \frac{1}{\Omega \sin \gamma} \arctan \left(\frac{f \circ g(c_1) - x_C}{g(c_1) - y_C - \frac{R_r}{\sin \gamma}} \right) \implies t = c_2 + h(c_1). \quad (74c)$$

By virtue of the (quite rather general) BCs (11), we obtain⁹

$$\begin{aligned} A_x(c_1, c_2) &= -\tilde{u}_x(f \circ g(c_1)) \cos(\varepsilon_\psi \omega_z c_2 + \varepsilon_\psi \omega h(c_1)) \\ &\quad + \tilde{u}_y(g(c_1)) \sin(\varepsilon_\psi \omega_z c_2 + \varepsilon_\psi \omega h(c_1)), \end{aligned} \quad (75a)$$

$$\begin{aligned} A_y(c_1, c_2) &= -\tilde{u}_x(f \circ g(c_1)) \sin(\varepsilon_\psi \omega_z c_2 + \varepsilon_\psi \omega h(c_1)) \\ &\quad - \tilde{u}_y(g(c_1)) \cos(\varepsilon_\psi \omega_z c_2 + \varepsilon_\psi \omega h(c_1)), \end{aligned} \quad (75b)$$

and coming back to the original arguments of $c_1(\mathbf{x})$ and $c_2(\mathbf{x}, t)$ yields

$$\begin{aligned} A_x(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) &= \Psi_x(\mathbf{x}) \cos \left(\varepsilon_\psi \omega_z t - \frac{\varepsilon_\psi \omega_z}{V_r} \Sigma(\mathbf{x}) \right) \\ &\quad - \Psi_y(\mathbf{x}) \sin \left(\varepsilon_\psi \omega_z t - \frac{\varepsilon_\psi \omega_z}{V_r} \Sigma(\mathbf{x}) \right), \end{aligned} \quad (76a)$$

$$\begin{aligned} A_y(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) &= \Psi_x(\mathbf{x}) \sin \left(\varepsilon_\psi \omega_z t - \frac{\varepsilon_\psi \omega_z}{V_r} \Sigma(\mathbf{x}) \right) \\ &\quad + \Psi_y(\mathbf{x}) \cos \left(\varepsilon_\psi \omega_z t - \frac{\varepsilon_\psi \omega_z}{V_r} \Sigma(\mathbf{x}) \right), \end{aligned} \quad (76b)$$

where we have introduced

$$\Sigma(\mathbf{x}) \triangleq V_r(h(\mathbf{x}) - h \circ c_1(\mathbf{x})), \quad (77a)$$

$$\Psi_x(\mathbf{x}) \triangleq -\tilde{u}_x(f \circ g \circ c_1(\mathbf{x})), \quad (77b)$$

$$\Psi_y(\mathbf{x}) \triangleq -\tilde{u}_y(g \circ c_1(\mathbf{x})). \quad (77c)$$

Inserting the above expressions (76) into (66), replacing the time variable t with the travelled distance s , the speeds with the slip parameters, and resorting to elementary trigonometric relations, we obtain the final solution for the steady-state deflections of the

bristle

$$u_x^-(\mathbf{x}) = \Psi_x(\mathbf{x}) \cos\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma(\mathbf{x})\right) - \Psi_y(\mathbf{x}) \sin\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma(\mathbf{x})\right) + \tilde{u}_x(x), \quad (78a)$$

$$u_y^-(\mathbf{x}) = \Psi_x(\mathbf{x}) \sin\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma(\mathbf{x})\right) + \Psi_y(\mathbf{x}) \cos\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma(\mathbf{x})\right) + \tilde{u}_y(y), \quad (78b)$$

which are clearly independent of the time variable t (or, equivalently, s). The procedure outlined so far is rather abstract; to give a practical example of how the solutions look like, we can derive an analytical expression in the case of a rectangular contact patch. It is worth pointing out that, for a rectangular shape, the function $\Sigma(\mathbf{x})$ reads exactly as in (49), depending on the specific BC. In particular, since the BCs and the characteristic curves for the 2DCM are the same as for the 2DM, it is still possible to identify the same regions \mathcal{P}_1 , \mathcal{P}_2 and \mathcal{P}_3 of the contact patch in which the steady-state deflection of the bristle $\mathbf{u}_t^-(\mathbf{x})$ will be given by a different expression. Moreover, because of the BC prescription, the solutions in the subdomains \mathcal{P}_1 and \mathcal{P}_2 will be again continuous at the interface $\mathcal{C}_2(\mathbf{x})$, whilst a discontinuity will occur on $\mathcal{C}_3(\mathbf{x})$. Summarising, we will have three expressions for $\Sigma(\mathbf{x})$, namely $\Sigma_1(\mathbf{x})$, $\Sigma_2(\mathbf{x})$, $\Sigma_3(\mathbf{x})$, and three expressions for $\Psi_x(\mathbf{x})$ and $\Psi_y(\mathbf{x})$, respectively. More specifically, we can define

$$\Psi_{1x}(\mathbf{x}) \triangleq \frac{\sigma_y R_r}{\varepsilon_\psi \varphi} + \frac{l}{2} - x_C, \quad (79a)$$

$$\Psi_{2x}(\mathbf{x}) \triangleq \frac{\sigma_y R_r}{\varepsilon_\psi \varphi} + \sqrt{\Gamma(\mathbf{x}) - R_1^2}, \quad (79b)$$

$$\Psi_{3x}(\mathbf{x}) \triangleq \frac{\sigma_y R_r}{\varepsilon_\psi \varphi} - \sqrt{\Gamma(\mathbf{x}) - R_3^2}, \quad (79c)$$

$$\Psi_{1y}(\mathbf{x}) \triangleq -\frac{\sigma_x R_r}{\varepsilon_\psi \varphi} - \sqrt{\Gamma(\mathbf{x}) - R_0^2} \text{sign}(\gamma) + \frac{R_r}{\varepsilon_\psi \sin \gamma}, \quad (79d)$$

$$\Psi_{2y}(\mathbf{x}) \triangleq -\frac{\sigma_x R_r}{\varepsilon_\psi \varphi} + \frac{w}{2} \text{sign}(\gamma) - y_C + \frac{R_r}{\varepsilon_\psi \varphi}, \quad (79e)$$

$$\Psi_{3y}(\mathbf{x}) \triangleq -\frac{\sigma_x R_r}{\varepsilon_\psi \varphi} - \frac{w}{2} \text{sign}(\gamma) - y_C + \frac{R_r}{\varepsilon_\psi \varphi}, \quad (79f)$$

so that we have the following expressions for the longitudinal and lateral components of the steady-state bristle displacement in the three subdomains

$$u_{1x}^-(\mathbf{x}) = \Psi_{1x}(\mathbf{x}) \cos\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma_1(\mathbf{x})\right) - \Psi_{1y}(\mathbf{x}) \sin\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma_1(\mathbf{x})\right) + \tilde{u}_x(x), \quad \mathbf{x} \in \mathcal{P}_1^-, s \in \mathbb{R}_{\geq 0}, \quad (80a)$$

$$u_{2x}^-(\mathbf{x}) = \Psi_{2x}(\mathbf{x}) \cos\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma_2(\mathbf{x})\right) - \Psi_{2y}(\mathbf{x}) \sin\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma_2(\mathbf{x})\right) + \tilde{u}_x(x), \quad \mathbf{x} \in \mathcal{P}_2^-, s \in \mathbb{R}_{\geq 0}, \quad (80b)$$

$$u_{3x}^-(\mathbf{x}) = \Psi_{3x}(\mathbf{x}) \cos\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma_3(\mathbf{x})\right) - \Psi_{3y}(\mathbf{x}) \sin\left(\frac{\varepsilon_\psi \varphi}{R_r} \Sigma_3(\mathbf{x})\right) + \tilde{u}_x(x), \quad \mathbf{x} \in \mathcal{P}_3^-, s \in \mathbb{R}_{\geq 0}, \quad (80c)$$

$$u_{1y}^-(\mathbf{x}) = \Psi_{1x}(\mathbf{x}) \sin\left(\frac{\varepsilon_{\psi}\varphi}{R_r}\Sigma_1(\mathbf{x})\right) + \Psi_{1y}(\mathbf{x}) \cos\left(\frac{\varepsilon_{\psi}\varphi}{R_r}\Sigma_1(\mathbf{x})\right) + \tilde{u}_y(y), \quad \mathbf{x} \in \mathcal{P}_1^-, s \in \mathbb{R}_{\geq 0}, \quad (80d)$$

$$u_{2y}^-(\mathbf{x}) = \Psi_{2x}(\mathbf{x}) \sin\left(\frac{\varepsilon_{\psi}\varphi}{R_r}\Sigma_2(\mathbf{x})\right) + \Psi_{2y}(\mathbf{x}) \cos\left(\frac{\varepsilon_{\psi}\varphi}{R_r}\Sigma_2(\mathbf{x})\right) + \tilde{u}_y(y), \quad \mathbf{x} \in \mathcal{P}_2^-, s \in \mathbb{R}_{\geq 0}, \quad (80e)$$

$$u_{3y}^-(\mathbf{x}) = \Psi_{3x}(\mathbf{x}) \sin\left(\frac{\varepsilon_{\psi}\varphi}{R_r}\Sigma_3(\mathbf{x})\right) + \Psi_{3y}(\mathbf{x}) \cos\left(\frac{\varepsilon_{\psi}\varphi}{R_r}\Sigma_3(\mathbf{x})\right) + \tilde{u}_y(y), \quad \mathbf{x} \in \mathcal{P}_3^-, s \in \mathbb{R}_{\geq 0}. \quad (80f)$$

On the other hand, to derive the transient solution $\mathbf{u}_t^+(\mathbf{x}, t)$, we can proceed analogously as in the previous section, to get

$$\begin{aligned} u_x^+(\mathbf{x}, s) &= \left[u_{x0}(\mathbf{x}_0(\mathbf{x}, s)) - \tilde{u}_x(x_0(\mathbf{x}, s)) \right] \cos\left(\frac{\varepsilon_{\psi}\varphi}{R_r}s\right) \\ &\quad - \left[u_{y0}(\mathbf{x}_0(\mathbf{x}, s)) - \tilde{u}_y(y_0(\mathbf{x}, s)) \right] \sin\left(\frac{\varepsilon_{\psi}\varphi}{R_r}s\right) + \tilde{u}_x(x), \quad \mathbf{x} \in \mathcal{P}^+, s \in \mathbb{R}_{\geq 0}, \end{aligned} \quad (81a)$$

$$\begin{aligned} u_y^+(\mathbf{x}, s) &= \left[u_{x0}(\mathbf{x}_0(\mathbf{x}, s)) - \tilde{u}_x(x_0(\mathbf{x}, s)) \right] \sin\left(\frac{\varepsilon_{\psi}\varphi}{R_r}s\right) \\ &\quad + \left[u_{y0}(\mathbf{x}_0(\mathbf{x}, s)) - \tilde{u}_y(y_0(\mathbf{x}, s)) \right] \cos\left(\frac{\varepsilon_{\psi}\varphi}{R_r}s\right) + \tilde{u}_y(y), \quad \mathbf{x} \in \mathcal{P}^+, s \in \mathbb{R}_{\geq 0}, \end{aligned} \quad (81b)$$

in which the ID $\mathbf{x}_0(s)$ is exactly as in (54). Also in this case, albeit the transient solution is not affected by the partition of the contact patch in a formal way, it must be distinguished amongst a different expression depending on each subdomains of \mathcal{P} . Furthermore, we particularly emphasise that the relation $\mathbf{u}_t^-(\mathbf{x}) = \mathbf{u}_t^+(\mathbf{x}, \Sigma(\mathbf{x}))$ holds between the steady-state and transient solution in each subdomain, i.e. the two solutions are continuous at the interface between their respective domains \mathcal{P}^- and \mathcal{P}^+ , which is again represented by the travelling function $s = \Sigma(\mathbf{x})$. From a physical perspective, this is a direct consequence of the fact that the bristles always enter the contact patch undeformed.

Finally, to prove the consistency of the proposed solutions, we should verify that they are equivalent to the ones obtained in Section 3 for small values of the parameter ε_{ψ} . The calculations are, however, quite cumbersome, so we do not show the steps. Generally speaking, both the steady-state expressions and the transient one converge to the corresponding formulae found by means of the 2DM by expanding the trigonometric functions into Taylor series up to the first order and then taking the limit for $\varepsilon_{\psi} \rightarrow 0$. By transitivity, it can be also concluded that the 2DCM reduces to the 1DM for sufficiently small values of the coefficient ε_{ψ} and camber angle γ .

Remark 5.1: In the preceding analyses, with reference to each model, we assumed $\mathbf{u}_{t0}(\mathbf{x}) \in C^1(\overset{\circ}{\mathcal{P}})$. However, if the slip inputs are varied before the transient corresponding to a previous condition is extinguished, the initial conditions are automatically only $C^0(\overset{\circ}{\mathcal{P}})$. In this case, the analytic solutions obtained by applying the method of the characteristics

obviously do not solve the original PDEs, but still represent the only candidate solutions for the weak formulation of the problem. For examples, the transient results shown in [52] for nonzero initial condition clearly represent a weak solution of the IDM (a strong solution is, indeed, not possible).

6. Model comparison

In this section, we compare the model developed so far against each other and with the classic IDM. The preceding analyses highlighted different microscopic phenomena which are not captured by the IDM; macroscopically, these discrepancies may be reflected in a different relation between the generalised forces acting on the tyre and the translational and rotational slips, e.g. the $F_x - \sigma_x$ curve.

With reference to the kinematic quantities (basically the slip inputs and the other parameters introduced in 2.1.1), it is firstly necessary to select a minimum set of independent coordinates which fully describe the problem. We can generally assume the existence of a hypersurface Υ so that

$$\Upsilon \left(\mathbf{F}_t, \dot{\mathbf{F}}_t, M_z, \dot{M}_z, \boldsymbol{\sigma}, \varphi, \varepsilon_\psi \right) = 0, \quad (82)$$

and invoking the IFT we can explicit the variables in interest as a function of the other ones. Perhaps, the most intuitive choice would be to choose the triad $(\boldsymbol{\sigma}, \varphi, \varepsilon_\psi)$.

To assess the performance of each model, we carried out different set of simulations in MATLAB[®] environment. For the case of a rectangular contact shape and under the assumption of vanishing sliding ($\mu_s \rightarrow \infty$), the total transient shear stress $q_t(\mathbf{x}, s) = \|\mathbf{q}_t(\mathbf{x}, s)\|$ acting in the contact patch is shown in Figure 5 for the three different models 1DCM, 2DM and 2DCM and for two values of the normalised travelled distance $\bar{s} = 1/2$ and 2, respectively. The pictures refer to the following kinematic parameters: $\sigma_x = \sigma_y = 0.1$, $\varphi = 1$ (corresponding to a value of the camber angle of $\gamma = 30^\circ$), $\varepsilon_\psi = 0.5$. Generally speaking, it is possible to note an overall agreement in the trend of the predicted shear stress; in particular, the values obtained by employing the 1DCM are slightly higher due to the fact that the bristle displacements are constrained to be zero only at $x = l/2$.

The assumption of vanishing sliding was then removed to investigate the steady-state problem. Again in the case of a rectangular contact shape, we assumed a parabolic pressure distribution of the type

$$q_z(\mathbf{x}) = \frac{6F_z}{wl} \left(\frac{l^2}{4} - x^2 \right), \quad (83)$$

where F_z is the total vertical force acting on the tyre. We also approximated the longitudinal and lateral components of the micro-sliding speed of the bristles in Equation (1b) as follows:

$$v_{sx}(y) \approx -V_r \left[\sigma_x - \frac{\varphi}{R_r} (y - y_C) \right], \quad (84a)$$

$$v_{sy}(x) \approx -V_r \left[\sigma_y + \frac{\varphi}{R_r} (x - x_C) \right], \quad (84b)$$

which is a common established approach when dealing with the IDM¹⁰. Note that the above simplification turns Equation (1b) from a PDE into an algebraic relation.

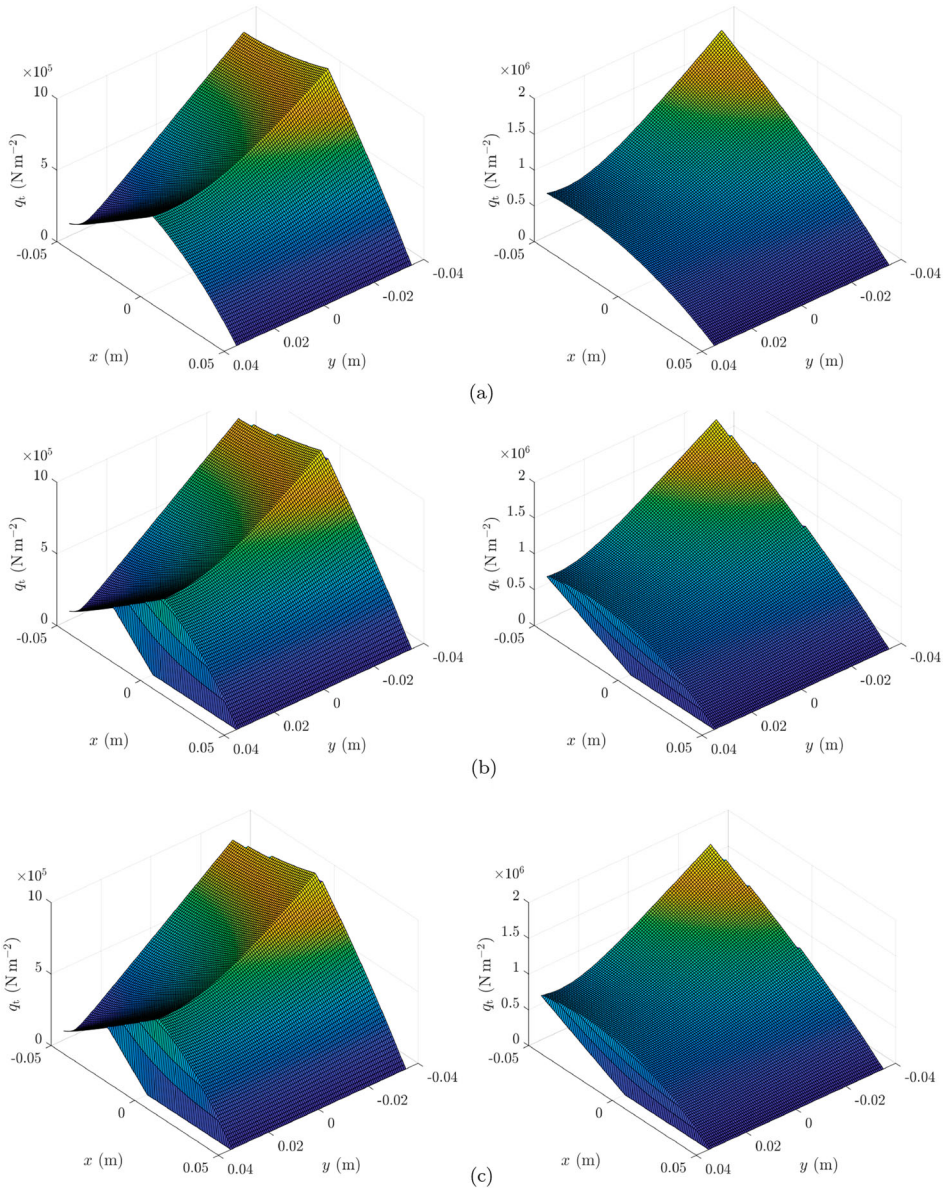


Figure 5. Transient trend of the total shear stress $q_t(\mathbf{x}, s)$ ($k_x = 8 \cdot 10^7 \text{ N m}^{-3}$, $k_y = 0.7k_x$) predicted by the different models for two values of the normalised travelled distance $\bar{s} = 1/2$ (left-hand side subplot) and $\bar{s} = 2$ (right-hand side subplot). The figures refer to the following values of the kinematic parameters: $\sigma_x = \sigma_y = 0.1$, $\varphi = 1$, $\varepsilon_\psi = 0.5$. (a) Total shear stress $q_t(\mathbf{x}, s)$ predicted using the 1DCM. The left and right-hand side subplots refer to a value of the normalised travelled distance of $\bar{s} = 1/2$ and 2, respectively. (b) Total shear stress $q_t(\mathbf{x}, s)$ predicted using the 2DM. The left and right-hand side subplots refer to a value of the normalised travelled distance of $\bar{s} = 1/2$ and 2, respectively. (c) Total shear stress $q_t(\mathbf{x}, s)$ predicted using the 2DCM. The left and right-hand side subplots refer to a value of the normalised travelled distance of $\bar{s} = 1/2$ and 2 respectively.

For each model under consideration, the steady-state tyre characteristics were computed numerically starting from the relations¹¹

$$\mathbf{F}_t(\boldsymbol{\sigma}, \varphi, \varepsilon_\psi) = \iint_{\mathcal{P}} \mathbf{q}_t(\mathbf{x}; \boldsymbol{\sigma}, \varphi, \varepsilon_\psi) \, dx \, dy, \quad (85)$$

$$M_z(\boldsymbol{\sigma}, \varphi, \varepsilon_\psi) = \iint_{\mathcal{P}} (\mathbf{x} - \mathbf{x}_C + \mathbf{u}_t(\mathbf{x}; \boldsymbol{\sigma}, \varphi, \varepsilon_\psi)) \times \mathbf{q}_t(\mathbf{x}; \boldsymbol{\sigma}, \varphi, \varepsilon_\psi) \, dx \, dy. \quad (86)$$

6.1. Analysis for constant geometric parameters

This analysis was conducted under the assumption of fixed geometric parameters, namely R_r , x_C and y_C . Furthermore, since all the models are supposedly equivalent for sufficiently small values of φ and ε_ψ , we only investigated the cases in which this two parameters were meant to have a significant impact over the tyre characteristics. Hence, we fixed $\varphi = 1$ and considered the two cases $\varepsilon_\psi = 0.1$ and $\varepsilon_\psi = 0.9$. The first condition describes a scenario in which a heavily cambered tyre ($\gamma \approx 50^\circ$) tyre is subjected to a low steering speed; the latter, instead, depicts a situation in which the camber angle is limited and the steering speed represents the main component of the spin variable.

The tyre was assumed anisotropic with $k_x = 8 \cdot 10^7 \, \text{N m}^{-3}$ and $k_y = 0.7 \cdot k_x$. The total vertical force on the tyre was set to $F_z = 4000 \, \text{N}$, whilst the static and dynamic friction coefficient were chosen to be $\mu_s = 0.9$ and $\mu_d = 0.7$. The coordinates of the wheel centre were assumed to be both zero, i.e. $x_C = y_C = 0$. For the rolling radius we chose $R_r = 0.28$ and we assigned the contact patch length and width as $l = 0.1$ and $w = 0.07 \, \text{m}$, respectively.

The longitudinal and lateral tyre characteristics are shown in Figure 6 versus the longitudinal slip σ_x and for different discrete values of the lateral one σ_y . Regarding the forces acting on the tyre, it is possible to note that high values of the camber angle ($\varepsilon_\psi = 0.1$) or the steering speed ($\varepsilon_\psi = 0.9$) have a negligible impact on the steady-state values. Variations in the trend of the self-aligning moment M_z , depicted in Figure 7, are instead appreciable in both cases; this is probably due to the fact that, for the computation of the moment, the exact distribution of the tangential stresses inside the contact patch plays a crucial role. More specifically, in the first scenario ($\varepsilon_\psi = 0.1$), the 2DM and 2DCM predict a similar asymmetric trend for each value of σ_y , whereas the 1DM and 1DCM are not able to capture this phenomenon. Of course, the solution provided by the 2DM aligns quite well with the one found by employing the 2DCM, since the values of the steering ratio are sufficiently small to disregard the coupling between the slips and spin parameter.

On the contrary, for $\varepsilon_\psi = 0.9$, the 1DM and 2DM fail in predicting the symmetric trend of the self-aligning moment, whilst the theories for coupled slips and spin both predict a similar trend, with the values foreseen by the 2DCM being slightly lower. From a physical perspective, this can be explained intuitively by considering that the two-dimensional theories prescribe multiple BCs on the edge of the rectangular domain. Since the BCs impose zero deflection of the bristles inflowing the contact patch, this means that they are constrained to adhere to the road in a wider area, and the growth rate in the spatial dimension is more limited.

In Figure 8, the relations $F_x - F_y$ are depicted again for different discrete values of σ_y and for the two cases $\varepsilon_\psi = 0.1$ and 0.9 , respectively. This kind of plot is very useful and it is often referred to *friction ellipse*, since all the point fall within an ellipse which gives the

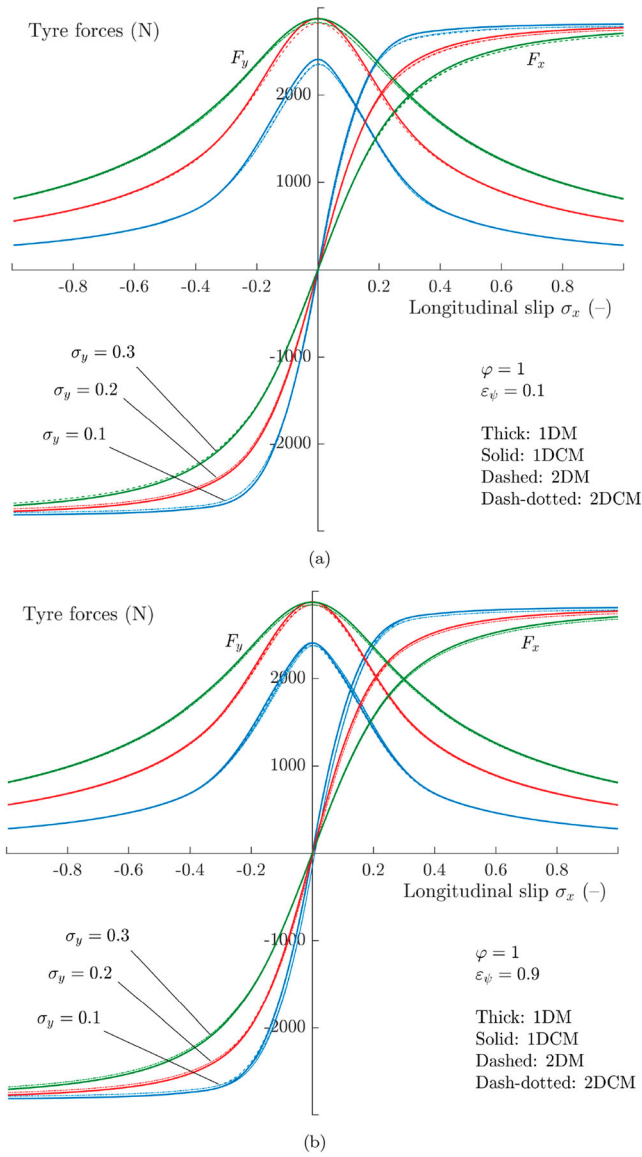


Figure 6. Tyre characteristics versus the longitudinal slip σ_x for different discrete values of the lateral slip σ_y and steering ratio ε_ψ . (a) Longitudinal and lateral tyre characteristics versus the longitudinal slip σ_x for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.1$. (b) Longitudinal and lateral tyre characteristics versus the longitudinal slip σ_x for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.9$.

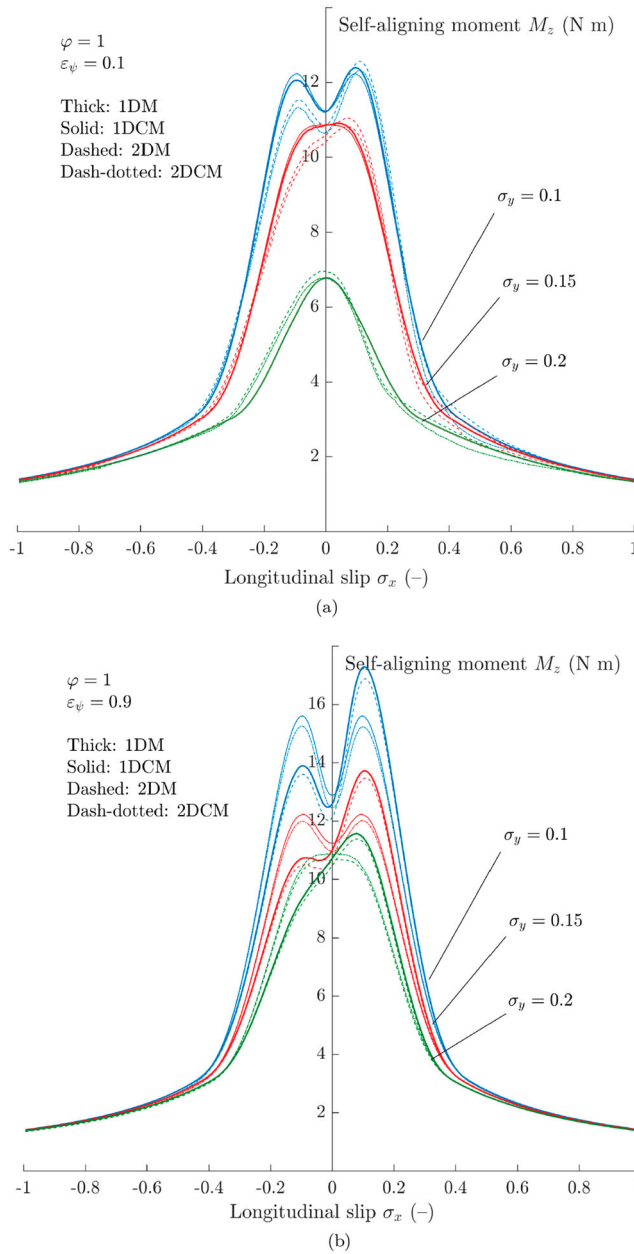


Figure 7. Self-aligning moment versus the longitudinal slip σ_x for different discrete values of the lateral slip σ_y and steering ratio ε_ψ . It can be noted that, for small values of the steering ratio, the 2DM and 2DCM both succeed in estimating the true trend of the self-aligning moment, where the other theories fail; conversely, for larger values of ε_ψ , the trend is better predicted by the 1DCM and 2DCM. (a) Self-aligning moment versus the longitudinal slip σ_x for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.1$. (b) Self-aligning moment versus the longitudinal slip σ_x for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.9$.

maximum value of the tangential force which can be exerted on the tyre itself. The previous considerations can be extended automatically to the analysis of Figure 8: for small values of the steering ratio, the 2DM and the 2DCM are totally equivalent and predict the same quantitative trend; on the other hand, as the parameter ε_ψ increases, the 2DCM aligns more with the 1DCM, since the coupling between the slips and the spin variable becomes preponderant.

A last comparison is performed with reference to the $F_y - M_z$ diagram in Figure 9, where the relations are plotted at constant values of the longitudinal slip σ_x . The conclusions which can be inferred are the same as for the previous analyses: the two-dimensional models do not exhibit any discrepancy for $\varepsilon_\psi = 0.1$, whilst at larger values of ε_ψ the mismatch increases and the 1DCM converges to the 2DCM.

Generally speaking, the novel models do not predict significant differences from the already-known tyre characteristics from the 1DM. Indeed, the maximum values of the tangential stresses acting in the contact patch is always limited by the available friction; if the same pressure distribution is assumed for all the models, and the structural and geometrical parameters are kept constant, it is reasonable to conjecture that the overall result would not change meaningfully.

6.2. Analysis for varying geometric parameters

High-cambered tyres are expected to undergo significant variations in the geometrical parameters. More specifically, a simple approximation which can be used to model the rolling radius as a function of the camber angle is $R_r(\gamma) \approx R_\delta \cos^2 \gamma$, where R_δ is the deformed radius of the tyre due to a pure vertical load. Accordingly, the lateral coordinate of the wheel hub centre becomes $y_C(\gamma) \approx R_r \sin \gamma / \cos^2 \gamma$, whilst the longitudinal deflection can be assumed to be almost zero. A second theoretical investigation was hence aimed at understanding how these geometrical variations influence the results predicted by the different theories. In particular, to make a fair comparison, we introduced some modifications in the 1DM to account for the term y_C in the steady-state equations for the deflection of the bristle. In particular, we generalised the steady-state solution for the bristle displacements in (29) as

$$u_x^-(\xi) = \sigma_x \xi - \frac{\varphi}{R_r} \xi (\eta - y_C), \quad (87a)$$

$$u_y^-(\xi) = \sigma_y \xi + \frac{\varphi}{R_r} \xi \left(x_{\mathcal{L}} - \frac{\xi}{2} - x_C \right). \quad (87b)$$

Simulation results are illustrated in Figure 10 for different values of σ_y and $\varphi = 0.7$, $\varepsilon_\psi = 0.1$. Generally speaking, even for higher values of ε_ψ , the coupling between the slips and spin parameters is negligible when the tyre experiences high levels of cambering, and only the 2DM and 2DCM are able to predict substantial discrepancies from the classic theory.

7. Discussion and conclusion

Brush models are a basic but still effective approach to physical tyre modelling. Indeed, they describe tyre characteristics by means of a minimum set of parameters, which are

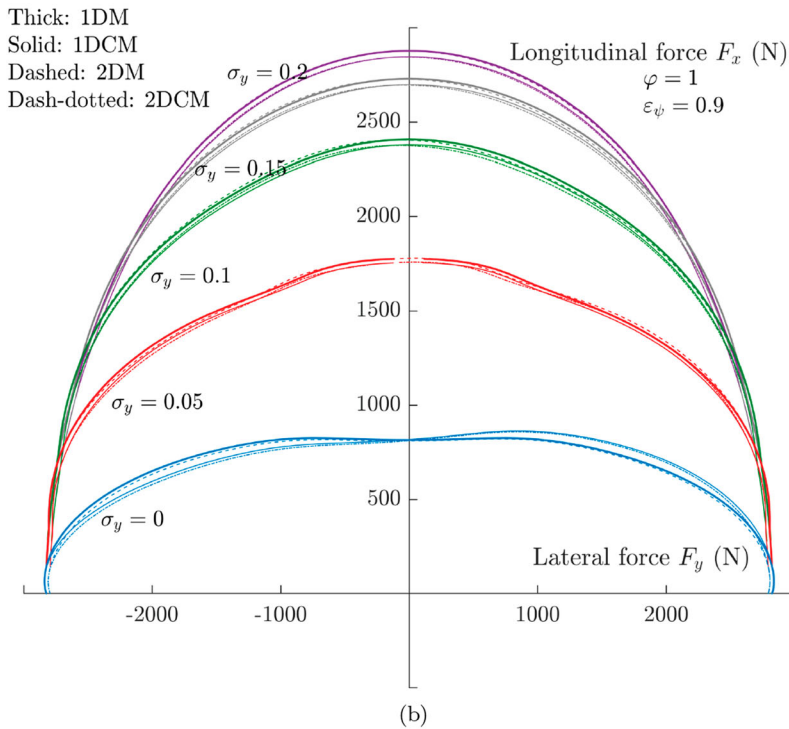
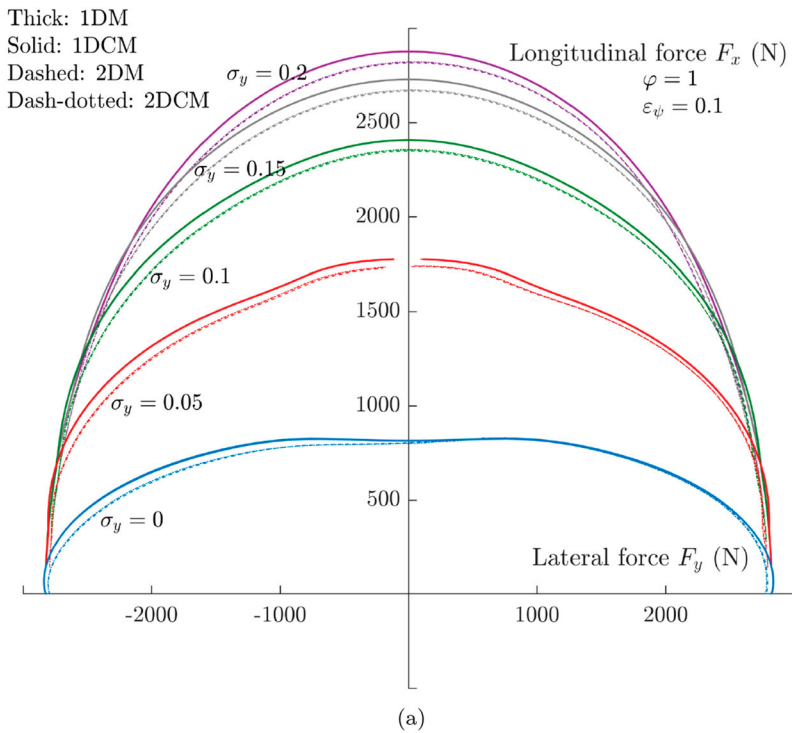


Figure 8. Friction ellipses for different discrete values of the lateral slip σ_y and steering ratio ε_{ψ} . It can be noted that, for small values of the steering ratio, the 2DM and 2DCM both succeed in estimating the true trend of the self-aligning moment, where the other theories fail; conversely, for larger values of ε_{ψ} , the trend is better predicted by the 1DCM and 2DCM. (a) Friction ellipse for different values of the lateral slip σ_y and steering ratio $\varepsilon_{\psi} = 0.1$. (b) Friction ellipse for different values of the lateral slip σ_y and steering ratio $\varepsilon_{\psi} = 0.9$.

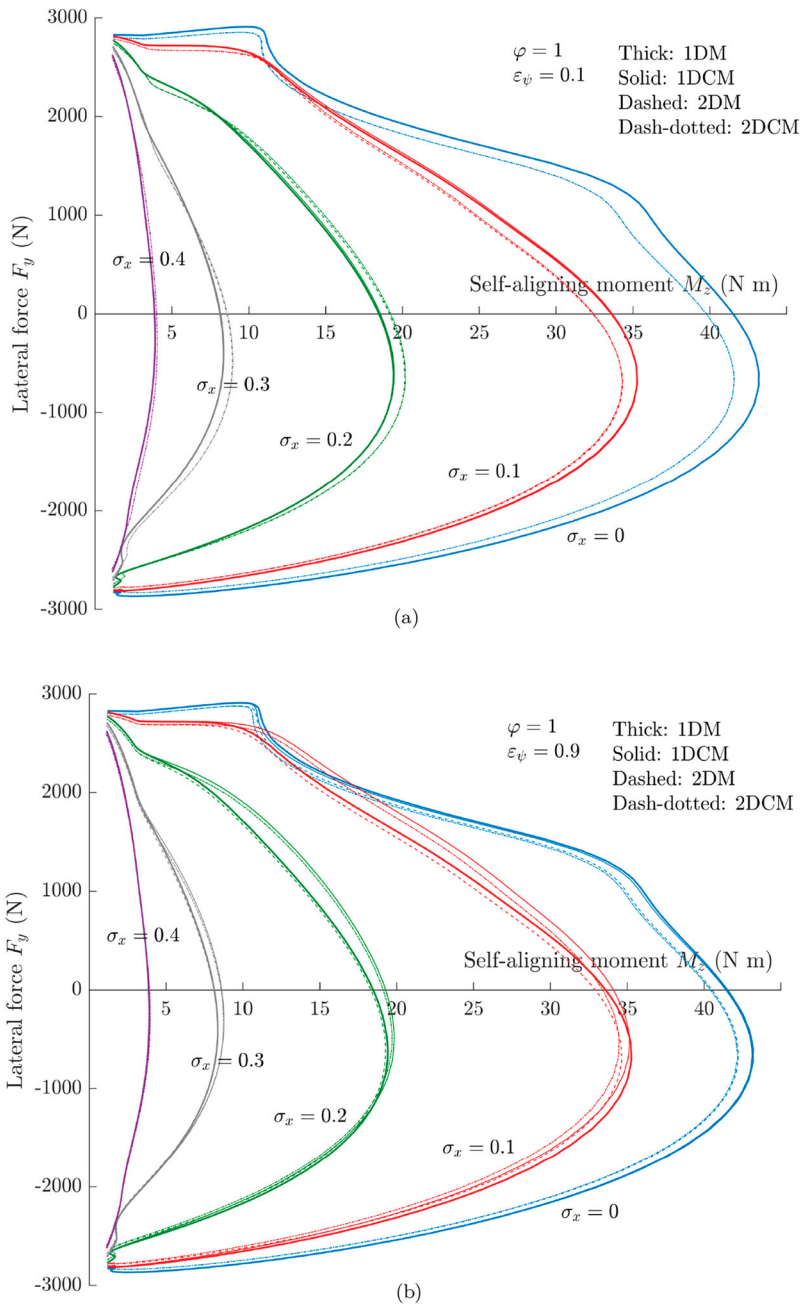


Figure 9. $F_y - M_z$ diagram for different discrete values of the lateral slip σ_y and steering ratio ε_ψ . It can be noted that, for small values of the steering ratio, the 2DM and 2DCM both succeed in estimating the true trend of the self-aligning moment, where the other theories fail; conversely, for larger values of ε_ψ , the trend is better predicted by the 1DCM and 2DCM. (a) $F_y - M_z$ diagram for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.1$. (b) $F_y - M_z$ diagram for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.9$.

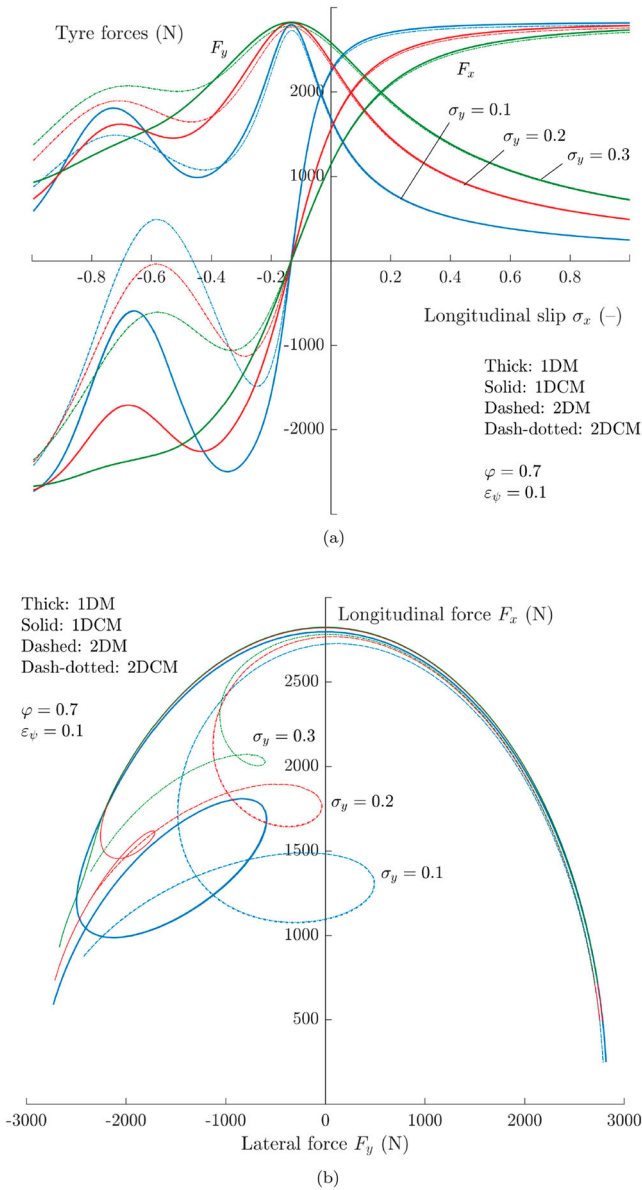


Figure 10. Tyre forces and friction ellipse for an heavily cambered tyre ($\gamma \approx 40^\circ$). When the tyre experiences high level of cambering, the results predicted by the two-dimensional theories exhibit appreciable differences with the ones found by means of the 1DM and 1DCM. (a) Longitudinal and lateral tyre characteristics versus the longitudinal slip σ_x for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.1$ and 0.9, respectively. The tyre rolling radius and the lateral coordinate of the wheel hub centre modelled as a function of the camber angle γ . (b) Friction ellipse for different values of the lateral slip σ_y and steering ratio $\varepsilon_\psi = 0.1$. The tyre rolling radius and the lateral coordinate of the wheel hub centre modelled as a function of the camber angle γ .

easy to interpret and to tune, and provide a quite exhaustive understanding of some interesting phenomena concerning the tyre-road interaction. Also, even though brush models are based on some rather simplifying assumptions, they are still able to predict measured forces.

In the present paper, we have extended the brush modelling beyond the standard theory to investigate the effect of the coupling between the slips and a two-dimensional velocity field inside the contact patch. More specifically, we have derived three different models of increasing order of complexity to analyse the contribution of larger spins and camber angles. The first theory concerns the presence of nonlinear relations between the slip parameters and is referred to as *one-dimensional model for coupled slips and spin* (1DCM). Both for the steady-state and the transient case, we have shown that, when the spin parameter and the steering ratio, φ and ε_{ψ} , respectively, are sufficiently large, the deflections of the bristle in the contact patch can be described by nonlinear functions depending on the space variables and the spin itself. The whole formulation is consistent with the classic brush theory, as confirmed by the asymptotic analysis which we have carried out for small values of the steering ratio ε_{ψ} and the spin φ .

The second variant which we have developed is the so-called *two-dimensional model* (2DM), and accounts for large values of the camber angle γ . When the camber is large, indeed, the total speed of a bristle may not be approximated by means of the *rolling speed* in longitudinal direction, and a more detailed description is needed. In this case, the modelling is further complicated by the fact that multiple BCs must be prescribed, yielding different form of the steady-state solution. The analysis carried out for a rectangular shape shows that it is possible to identify three regions inside the contact patch in which the solution for the vector displacement is provided by a different set of equations. These expressions are highly nonlinear with respect to the space coordinates, but no coupling is predicted between the slip variables. In each domain, the transient solution is again continuous at the transition with the steady-state one. The asymptotic analysis has shown that the novel theory is equivalent to the standard one for small values of camber angles.

Finally, the most general *two-dimensional model for coupled slips and spin* (2DCM) accounts for both the coupling between the slip parameters and for a two-dimensional velocity field. This model can thus be used when a cambered tyre experiences high steering speeds. The underlying mathematics which the theory is grounded on is, of course, more complicated. However, a procedure to derive the general solution in parametric form has been indicated and some qualitative considerations about its properties have been discussed. Indeed, the characteristics equations are the same as for the linear case, so that some conclusions can be easily extended to the nonlinear theory. In particular, we have shown that three expressions for the steady-state solution can be found in the same subdomains as for the linear model, and that these solutions are independent on the time variable. The explicit formulae provided for the steady-state and transient longitudinal deflections $\mathbf{u}_t^-(\mathbf{x}, s)$ and $\mathbf{u}_t^+(\mathbf{x}, s)$, respectively, share some similarities with the solutions found in the case of the two simpler theories developed in this paper. More specifically, the 2DCM theory predicts the coupling between the slip (in turn) and spin parameters and the dependency, for each component of the bristle displacement, on both the translational slip variables.

Limiting the attention to the brush theory, the new models presented in this paper could be useful to study some phenomena which are often neglected by the simplest classic model and deserve to be explored in greater detail. For example, the influence of large camber angles is preponderant during the cornering of motorcycles. Also, recent innovations in the motorsport scenario require extensive investigations about sudden variations in the camber setup. The theory could hence be extended towards real-time applications. Of course, the first step would be experimental validation. Advanced FEM or Multibody tyre models could serve this purpose, or alternatively data collected on a flat track testing machine. However, there are many other directions in which the present work can be potentially expanded. For example, we have already mentioned that the brush models, albeit being effective and particularly pleasing for pedagogical purposes, are probably a very naive approach to physical tyre modelling. As remarked by several authors, in the last years many efforts have been devoted to the development of accurate friction models which may represent a valuable addition to the models developed in the present paper. A natural extension of this work, for example, would be to combine the 2DM with the LuGre formulation to account for a variable velocity field inside the contact patch. It is interesting to note that, for example, an analytical solution to the unsteady-state LuGre model is rarely indicated in literature. The hybridisation with a friction model would also allow to overcome some limiting assumptions introduced in this theoretical analysis. Especially in the context of fast and accurate vehicle dynamics simulation, it is universally acknowledged that static or quasi-static brush models meet annoying difficulties in handling low-speed situations, whereas dynamic models perform much better. Usually, dynamic effects are related to the compliance of the tyre carcass, which has been systematically neglected in the present investigation.

Several strategies proposed in literature could hence be browse. One option could be to conduct a similar study for the *stretched-string tyre model* [2], which accounts for a distributed deformation of the tyre carcass along the longitudinal direction and has represented the starting point for several derivative models, e.g. the *single-point contact point model* [2]. However, the assumption of constant deformation in the lateral direction would eventually lead to a paradoxical result, and the incompatibility should be addressed by considering more complex models from advanced beam or membrane theories or by introducing some averaging methods. In this sense, it would be possible to follow, for example, Deur et al. [17] and provide an approximated expression for the function describing the carcass centreline. A recent solution based on the interpolation between the steady-state and transient solution of the tyre-road kinematic equations and called *two-regime tyre formulae* (TRF) has been developed, amongst the others, by the authors in [52]. This could represent a viable alternative to some stretched-string or LuGre-based enhanced formulations of the classic brush models.

Another possible choice would be to integrate the proposed theories with tyre models developed *ad hoc* for vehicle dynamics simulations. Whilst taking into account complex road geometries – e.g. uneven roads and soft soil – directly in the brush theory might be prohibitively challenging, it would definitely be beneficial to lean towards hybrid or semi-empirical solutions. To this extent, different algorithms have been developed to determine the position of the (lumped) contact point between the tyre and the road in case of irregular surfaces. These algorithms can be either based on numerical techniques [63] or simplifications of already-existing tyre models [64]. Combining a three-dimensional

wheel model with a description of the generalised forces acting on the tyre borrowed from another theory would consist in a quite consolidated approach, and should be taken in consideration.

Notes

1. In this paper, we use *tangential* and *planar* interchangeably.
2. In general, $\mathbf{x} = \mathbf{x}(t)$ is a three-dimensional position vector $\mathbf{x}(t) = (x(t), y(t), z(t))$. However, the tangential deflection of a bristle $\mathbf{u}_t(\mathbf{x}, t)$ is always evaluated at $z = 0$, so we write, with some abuse of notation, $\mathbf{u}_t(\mathbf{x}, t) = \mathbf{u}_t(x, y, t)$ in extended form instead of $\mathbf{u}_t(x, y, 0, t)$. Equivalently, if the travelled distance s is used as independent variable, we write $\mathbf{u}_t(\mathbf{x}, s) = \mathbf{u}_t(x, y, t)$ instead of $\mathbf{u}_t(x, y, 0, s)$.
3. The subscript s stands for *sliding*.
4. We discriminate between *macro* parameters, which are independent of the vector position \mathbf{x} , and *micro* parameters, which depend on \mathbf{x} and are defined locally.
5. Of course, every quantity can be a parameter or a variable depending on the specific problem at hand.
6. The main reason for assuming no sliding is that, in case of finite friction, the IC should be prescribed on an open set which would be unknown *a priori*. We could also think of defining the IC on the whole interior \mathcal{P} of the contact patch and then construct the solution to satisfy Equation (1b), but this would be formally incorrect.
7. Note that $G_x(\cdot, \cdot)$ and $G_y(\cdot, \cdot)$ are formally identical.
8. The function $f(\cdot)$ would correspond to $f_i(\cdot) = x_{\mathcal{L}_i}(\cdot)$, $i = 1, \dots, I$, but we prefer to use $f(\cdot)$ to lighten the notation.
9. In this paper, the rule for the composition of two functions $f \circ g(\cdot)$ must be read as $f \circ g(\cdot) = f(g(\cdot))$.
10. In general, however, committing this approximation may lead to improper results.
11. We have used the semicolon to distinguish between space coordinates and slip variables. Note that, whilst for the deflection of the bristle the generalised slips can be interpreted as parameters, for the tyre characteristics they represent the main variables. Of course, everything depends on the problem at hand.

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Nomenclature

Forces and Moments	Unit	Description
q_t	N m^{-2}	Tangential shear stress vector
q_t	N m^{-2}	Total tangential shear stress
F_t	N	Tangential force vector
F_x	N	Longitudinal tyre force
F_y	N	Longitudinal tyre force
F_z	N	Vertical force acting on the tyre
M_z	N m	Self-aligning moment
Displacements	Unit	Description
u_t	m	Bristle tangential deflection vector
u_t^-	m	Steady-state tangential displacement vector of the bristle
u_t^+	m	Transient tangential displacement vector of the bristle
u_{t0}	m	Initial tangential displacement vector of the bristle (IC)
u_x^-	m	Steady-state longitudinal deflection
u_x^+	m	Transient longitudinal deflection
u_{x0}	m	Initial longitudinal deflection (IC)
u_y^-	m	Steady-state lateral deflection
u_y^+	m	Transient lateral deflection
u_{y0}	m	Initial lateral deflection (IC)
s	m	Travelled distance
x	m	Planar coordinate vector
x_0	m	Initial data vector (ID)
x	m	Longitudinal coordinate
x_0	m	Initial lateral datum (ID)
y	m	Lateral coordinate
y_0	m	Initial lateral datum (ID)
η	m	Alternative lateral coordinate
ξ	m	Alternative longitudinal coordinate or distance from the entrance
ζ	m	Alternative vertical coordinate
Speeds	Unit	Description
v	m s^{-1}	Three-dimensional velocity field
v_t	m s^{-1}	Tangential velocity field
v_x	m s^{-1}	Longitudinal component of the velocity field
v_y	m s^{-1}	Lateral component of the velocity field
v_{st}	m s^{-1}	Micro-sliding tangential speed vector
v_{sx}	m s^{-1}	Micro-sliding longitudinal speed
v_{sy}	m s^{-1}	Micro-sliding lateral speed
V_r	m s^{-1}	Tyre rolling speed
V_x	m s^{-1}	Longitudinal speed of the wheel hub
V_y	m s^{-1}	Lateral speed of the wheel hub
V_{st}	m s^{-1}	Tangential macro-sliding speed vector
V_{sx}	m s^{-1}	Longitudinal macro-sliding speed
V_{sy}	m s^{-1}	Lateral macro-sliding speed
$\dot{\psi}$	rad s^{-1}	Steering speed
ω_z	rad s^{-1}	Angular speed around the z axis
Ω	rad s^{-1}	Angular speed of the rim
Slip Parameters	Unit	Description
ε_γ	–	Camber ratio
ε_ψ	–	Steering ratio
σ	–	Translational slip vector
σ_x	–	Longitudinal slip
σ_y	–	Lateral slip
φ	–	Rotational slip or spin parameter

Geometrical Parameters	Unit	Description
l	m	Contact patch length
w	m	Contact patch width
\mathbf{x}_C	m	Wheel hub centre coordinate vector
x_C	m	Wheel hub centre longitudinal coordinate
$x_{\mathcal{L}}$	m	Leading edge position
y_C	m	Wheel hub centre lateral coordinate
γ	rad	Camber angle
R_0	m	Radius of the circumference C_0
R_1	m	Radius of the circumference C_1
R_2	m	Radius of the circumference C_2
R_3	m	Radius of the circumference C_3
R_4	m	Radius of the circumference C_4
R_5	m	Radius of the circumference C_5
R_r	m	Rolling radius
R_z	m	Vertical radius
R_δ	m	Tyre deformed radius
Stiffnesses and Compliances	Unit	Description
k_x	N m^{-3}	Bristle longitudinal stiffness
k_y	N m^{-3}	Bristle lateral stiffness
\mathbf{K}_t	N m^{-3}	Matrix of the bristle tangential stiffnesses
Friction Parameters	Unit	Description
μ_d	–	Sliding friction coefficient
μ_s	–	Sticking friction coefficient
Functions and Operators	Unit	Description
D_t	s^{-1}	Tangential partial differential operator
∇_t	m^{-1}	Tangential gradient
$\Phi_X(\cdot)$	$\text{m}^3 \text{s}^{-1}$	Flux through X
Γ	m^2	Gamma function
Ξ	–	Xi function
$\Sigma_1, \Sigma_2, \Sigma_3$	m	Sigma functions
$\Phi_{1x}, \Phi_{2x}, \Phi_{3x}$	m	Phi functions for the longitudinal deflection
$\Phi_{1y}, \Phi_{2y}, \Phi_{3y}$	m	Phi functions for the lateral deflection
$\Psi_{1x}, \Psi_{2x}, \Psi_{3x}$	m	Psi functions for the longitudinal deflection
$\Psi_{1y}, \Psi_{2y}, \Psi_{3y}$	m	Psi functions for the lateral deflection
Sets	Unit	Description
\mathcal{P}	–	Contact patch
$\overset{\circ}{\mathcal{P}}$	–	Interior of \mathcal{P}
$\partial\mathcal{P}$	–	Boundary of \mathcal{P}
\mathcal{L}	–	Leading edge
\mathcal{N}	–	Neutral edge
\mathcal{T}	–	Trailing edge
$\mathbb{R}_{\geq 0}$	–	Set of positive real numbers (including 0)
$\mathbb{R}_{> 0}$	–	Set of strictly positive real numbers (excluding 0)

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Appendices

Appendix 1. Derivation of the steady-state solution for the two-dimensional theory

In this appendix we illustrate the procedure to derive the three expressions for the longitudinal displacement in the domains \mathcal{P}_1 , \mathcal{P}_2 , \mathcal{P}_3 . The same steps can be extended automatically to the lateral problem, so we only consider the longitudinal deflection.

From the first BC in (35) we get

$$G_x \left(c_1 \left(\frac{l}{2}, y \right), c_2 \left(\frac{l}{2}, y, t \right) \right) = -\sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) + \frac{\varphi}{\sin \gamma} \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right) + \frac{l}{2} \right], \quad (\text{A1})$$

where the arguments of $G_x(\cdot, \cdot)$ read

$$c_1 \left(\frac{l}{2}, y, t \right) = \left(\frac{l}{2} - x_C \right)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2, \quad (\text{A2a})$$

$$c_2 \left(\frac{l}{2}, y, t \right) = t - \frac{1}{\Omega \sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{y - y_C - \frac{R_r}{\sin \gamma}} \right). \quad (\text{A2b})$$

The above Equation (A2) can be solved for y and t to obtain

$$y = y_C + \frac{R_r}{\sin \gamma} \pm \sqrt{c_1 - \left(\frac{l}{2} - x_C \right)^2}, \quad (\text{A3a})$$

$$t = c_2 \pm \frac{1}{\Omega \sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{\sqrt{c_1 - \left(\frac{l}{2} - x_C \right)^2}} \right). \quad (\text{A3b})$$

Now, some physical considerations are needed to choose the correct sign in the above expressions. As we mentioned before, the characteristics equations for the problem are represented by circumferences of arbitrary radius centred in $C_\gamma = (x_C, y_C + \frac{R_r}{\sin \gamma})$ on the Oxy plane, which is the point where the wheel axis intercepts the road. This means that, if the camber angle γ is positive, this family of circumferences is centred on the right of the contact patch; conversely, on the left. Since the contact patch includes both points with negative and lateral coordinate, we have to choose the alternative with negative sign when solving for y . From the condition (32a), we have to take the solution with the negative sign for $\gamma > 0$; conversely, if $\gamma < 0$, recalling the second condition (32b), we need to take the solution with the positive sign. Thus, we get

$$G_x(c_1, c_2) = \sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{\sqrt{c_1 - \left(\frac{l}{2} - x_C \right)^2}} \right) \text{sign}(\gamma) - \frac{\varphi}{\sin \gamma} \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{\sqrt{c_1 - \left(\frac{l}{2} - x_C \right)^2}} \right) \text{sign}(\gamma) - \frac{l}{2} \right]. \quad (\text{A4})$$

Coming back to the original argument yields

$$\begin{aligned}
 & G_x(c_1(\mathbf{x}), c_2(\mathbf{x}, t)) \\
 &= \sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{\sqrt{(x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2 - \left(\frac{l}{2} - x_C\right)^2}} \right) \text{sign}(\gamma) \\
 &\quad - \frac{\varphi}{\sin \gamma} \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{\frac{l}{2} - x_C}{\sqrt{(x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2 - \left(\frac{l}{2} - x_C\right)^2}} \right) \text{sign}(\gamma) - \frac{l}{2} \right], \tag{A5}
 \end{aligned}$$

and combining the expression found above with (43a) yields the final formula for the steady-state longitudinal deflection $u_{1x}^-(\mathbf{x})$ in the region \mathcal{P}_1 , which can be restated in compact form by means of (47a), together with (46a), (49a) and (49d).

The second BC of (35) applies for $x > x_C$. With an analogous procedure as previously, we find

$$\begin{aligned}
 & G_x \left(c_1 \left(x, \frac{w}{2} \text{sign}(\gamma) \right), c_2 \left(x, \frac{w}{2} \text{sign}(\gamma), t \right) \right) \\
 &= \sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma}} \right) \\
 &\quad - \varphi \frac{1}{\sin \gamma} \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma}} \right) + x \right]. \tag{A6}
 \end{aligned}$$

Solving c_1 for x and t gives, respectively

$$x = x_C \pm \sqrt{c_1 - \left(\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma} \right)^2}, \tag{A7a}$$

$$t = c_2 \pm \frac{1}{\Omega \sin \gamma} \arctan \left(\frac{\sqrt{c_1 - \left(\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma} \right)^2}}{\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma}} \right), \tag{A7b}$$

where we have to choose the positive solution since the boundary is only defined for $x \geq x_C$. Substituting into (A6) and coming back to the original variables yields for the longitudinal deflection in the second domain \mathcal{P}_2 , reading as $u_{2x}^-(\mathbf{x})$ as in (47b), with $\Sigma_2(\mathbf{x})$ and $\Phi_{2x}(\mathbf{x})$ defined as in (49b)

and (49e). Finally, in the third domain \mathcal{P}_3 , the third BC applies, leading to

$$\begin{aligned}
 & G_x \left(c_1 \left(x, -\frac{w}{2} \text{sign}(\gamma) \right), c_2 \left(x, -\frac{w}{2} \text{sign}(\gamma), t \right) \right) \\
 &= \sigma_x \frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}} \right) \\
 & - \varphi \frac{1}{\sin \gamma} \left[\frac{R_r}{\sin \gamma} \arctan \left(\frac{x - x_C}{\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}} \right) - x \right]. \quad (\text{A8})
 \end{aligned}$$

As usual, solving c_1 and c_2 provides

$$x = x_C \pm \sqrt{c_1 - \left(\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma} \right)^2}, \quad (\text{A9a})$$

$$t = c_2 \mp \frac{1}{\Omega \sin \gamma} \arctan \left(\frac{\sqrt{c_1 - \left(\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma} \right)^2}}{\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}} \right), \quad (\text{A9b})$$

where we have to choose the negative solution for x (and the positive for the time t), since the BC applies only for $x < x_C$. We obtain the expression for $u_{3x}^-(\mathbf{x})$ as in (47c), with $\Sigma_3(\mathbf{x})$ and $\Phi_{3x}(\mathbf{x})$ defined as in (49c) and (49f), respectively.

Appendix 2. Determination of the domains for the steady-state solution

The subdomains of the contact patch for which we have the three different solutions must be determined starting from some considerations about the constraint on the space variables. For the solution $u_{2t}^-(\mathbf{x})$, we know, from Equation (A7a), that the following inequality must be verified:

$$x_C < x_C + \sqrt{(x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2 - \left(\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma} \right)^2} < \frac{l}{2}, \quad (\text{A10})$$

which constitutes the set of points between the two circumferences

$$\begin{aligned}
 \mathcal{C}_1(\mathbf{x}) &\triangleq \left\{ \mathbf{x} \in \mathbb{R}^3 \left| (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2 = R_1^2, z = 0 \right. \right\} \\
 &= \left\{ \mathbf{x} \in \mathbb{R}^3 \left| \Gamma(\mathbf{x}) = R_1^2, z = 0 \right. \right\}, \quad (\text{A11a})
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{C}_2(\mathbf{x}) &\triangleq \left\{ \mathbf{x} \in \mathbb{R}^3 \left| (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma} \right)^2 = R_2^2, z = 0 \right. \right\} \\
 &= \left\{ \mathbf{x} \in \mathbb{R}^3 \left| \Gamma(\mathbf{x}) = R_2^2, z = 0 \right. \right\}, \quad (\text{A11b})
 \end{aligned}$$

with

$$R_1 \triangleq \frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma}, \quad (\text{A12a})$$

$$R_2 \triangleq \sqrt{\left(\frac{w}{2} \text{sign}(\gamma) - y_C - \frac{R_r}{\sin \gamma}\right)^2 + \left(\frac{l}{2} - x_C\right)^2} = \sqrt{R_1^2 + R_0^2}, \quad (\text{A12b})$$

where R_0 reads as in (46a). Of course, we must only consider the points between the circles which also belong to the domain \mathcal{P} . Thus, the domain \mathcal{P}_2 can be written as

$$\mathcal{P}_2 \triangleq \left\{ \mathbf{x} \in \mathcal{P} \mid R_1^2 < (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2 < R_2^2 \right\} = \left\{ \mathbf{x} \in \mathcal{P} \mid R_1^2 < \Gamma(\mathbf{x}) < R_2^2 \right\}. \quad (\text{A13})$$

With similar reflections, to find the domain \mathcal{P}_3 , we note that from Equation (A9a) it must necessarily be

$$-\frac{l}{2} < x_C - \sqrt{(x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2} - \left(\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}\right)^2 < x_C, \quad (\text{A14})$$

which are the points contained by the implicit curves

$$\begin{aligned} \mathcal{C}_3(\mathbf{x}) &\triangleq \left\{ \mathbf{x} \in \mathbb{R}^3 \mid (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2 = R_3^2, z = 0 \right\} \\ &= \left\{ \mathbf{x} \in \mathbb{R}^3 \mid \Gamma(\mathbf{x}) = R_3^2, z = 0 \right\}, \end{aligned} \quad (\text{A15a})$$

$$\begin{aligned} \mathcal{C}_4(\mathbf{x}) &\triangleq \left\{ \mathbf{x} \in \mathbb{R}^3 \mid (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2 = R_4^2, z = 0 \right\} \\ &= \left\{ \mathbf{x} \in \mathbb{R}^3 \mid \Gamma(\mathbf{x}) = R_4^2, z = 0 \right\}, \end{aligned} \quad (\text{A15b})$$

with

$$R_3 \triangleq \frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}, \quad (\text{A16a})$$

$$R_4 \triangleq \sqrt{\left(\frac{w}{2} \text{sign}(\gamma) + y_C + \frac{R_r}{\sin \gamma}\right)^2 + \left(\frac{l}{2} + x_C\right)^2} = \sqrt{R_3^2 + R_5^2}, \quad (\text{A16b})$$

with R_5 defined as in (46f). The domain \mathcal{P}_3 is given by

$$\begin{aligned} \mathcal{P}_3 &\triangleq \left\{ \mathbf{x} \in \mathcal{P} \mid R_3^2 < (x - x_C)^2 + \left(y - y_C - \frac{R_r}{\sin \gamma}\right)^2 < R_4^2 \cap x < x_C \right\} \\ &= \left\{ \mathbf{x} \in \mathcal{P} \mid R_3^2 < \Gamma(\mathbf{x}) < R_4^2 \cap x < x_C \right\}. \end{aligned} \quad (\text{A17})$$

Finally, we can define the region \mathcal{P}_1 as the complementary of $\mathcal{P}_2 \cup \mathcal{P}_3$ to \mathcal{P} . We emphasise that the domain \mathcal{P}_1 is the only one predicted by the classic brush theory.